



A-Arevalo-Lopez  
C. minaud  
H. Nimoh



NEUTRONS  
FOR SCIENCE  
C .Colin



M. Lü



UNIVERSITÄT



R. Glaum

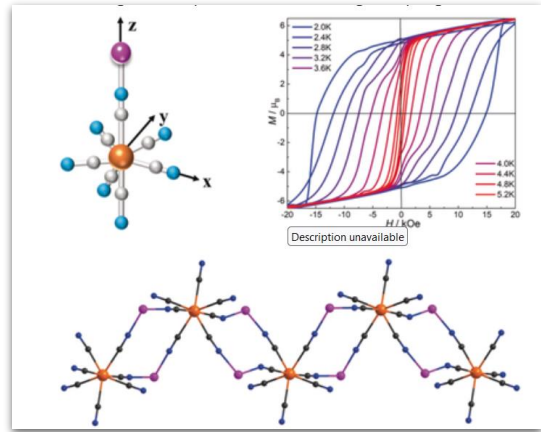
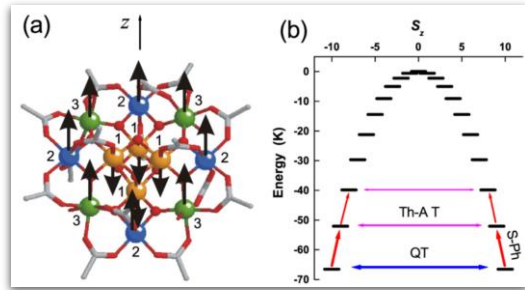


## SOME EXOTIC LOW-D MAGNETIC INORGANIC SYSTEMS, FRUSTRATION VS. MAGNETIC ORDERING

Olivier Mentré UCCS, Univ. Lille-Artois/Ecole Centrale/CNRS, France

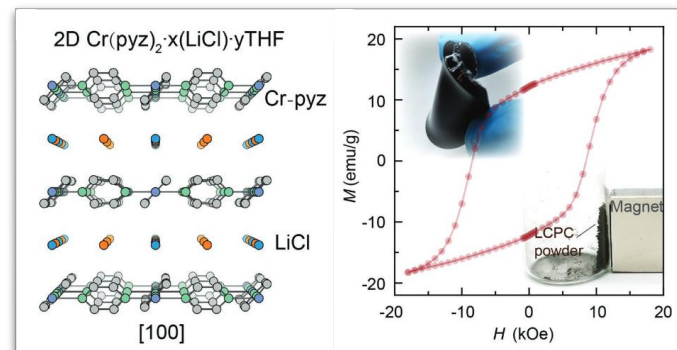
# Molecular magnetism

Single Molecule Magnet  
*quantum tunneling*



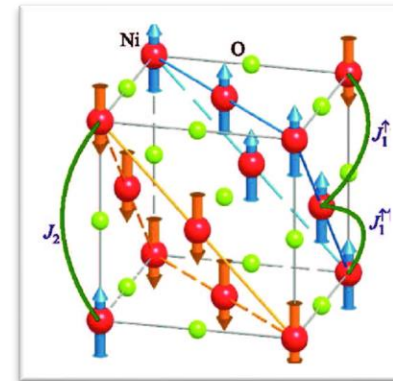
Single Chain Magnet  
*slow-spin relaxation → FM*

2D-Molecular magnets  
*anisotropy*



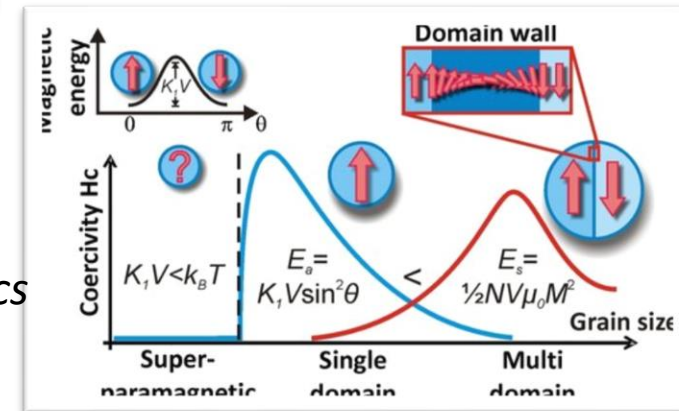
# Inorganic magnetic materials

Intermetallics  
*permanent magnets*



Magnetic Oxides, Sulfides, etc  
*3D magnetic ordering*

Nano-Magnetism  
*thin Films, particles, Spintronics*



# Exploration of Low-D Inorganic chemical systems

## Towards low-D units

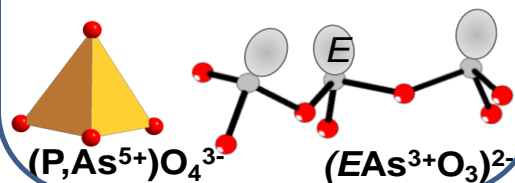


$r = 1.00 \text{ \AA}$

$1.18 \text{ \AA}$

$1.35 \text{ \AA}$

Alkali K, Rb, Cs



- Teflon-lined bomb : autogeneous pressure
- Hydrazine red. agent: stabilization low-redox

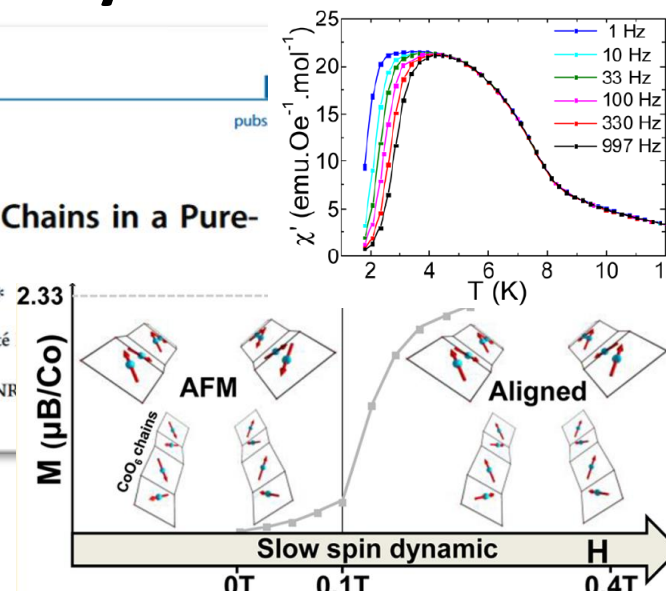
## Inorganic Chemistry

### Slow Spin Dynamics between Ferromagnetic Chains in a Pure-Inorganic Framework

Régnald David,<sup>†</sup> Houria Kabbour,<sup>†</sup> Silviu Colis,<sup>‡</sup> and Olivier Mentré<sup>†,\*</sup>

<sup>†</sup>UMR 8181 CNRS, Unité de Catalyse et de Chimie du Solide (UCCS USTL), Université d'Ascq, France

<sup>‡</sup>Institut de Physique et Chimie des Matériaux de Strasbourg (IPCMS), UMR 7504 CNRS (UDS-ECPM), F-67034 Strasbourg Cedex 2, France



## Inorganic Chemistry

Cite This: *Inorg. Chem.* 2019, 58, 12609–12617

Article  
pubs.acs.org/IC

### Metamagnetic Transitions versus Magnetocrystalline Anisotropy in Two Cobalt Arsenates with 1D Co<sup>2+</sup> Chains

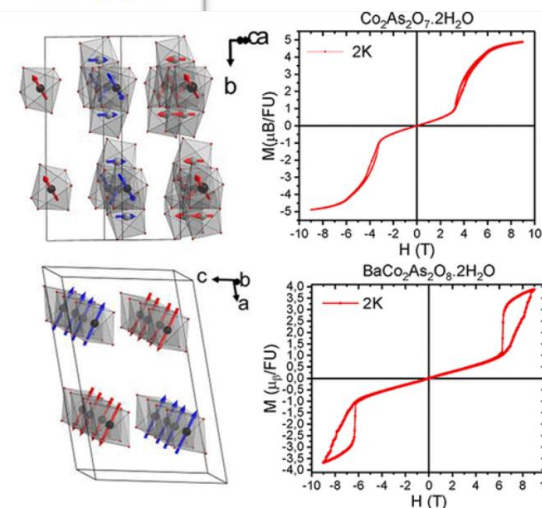
Bastien Leclercq,<sup>†</sup> Houria Kabbour,<sup>†</sup> Françoise Damay,<sup>‡</sup> Claire V. Colin, Angel M. Arevalo-Lopez,<sup>†</sup> and Olivier Mentré<sup>\*,†</sup>

<sup>†</sup>UCCS, UMR-CNRS 8181, Université Lille-ENSCL, Avenue Mendeleiev, 59655 Villeneuve

<sup>‡</sup>Laboratoire Léon Brillouin, UMR12-CNRS, CEA Saclay, Bât. 563, 91191 Gif-Sur-Yvette, Fr.

<sup>§</sup>Institut NEEL, CNRS/UGA UPR2940, CNRS/UGA UPR2940, 25 rue des Martyrs, 38042

<sup>||</sup>Laboratoire CRISTMAT, UMR 6508-CNRS, ENSICAEN, 6 Bd. Du Maréchal Juin, 14000



**aim of this « tutorial »** ... as requested by G. Chastanet

...

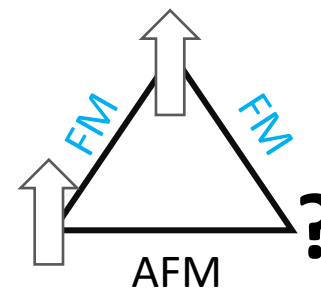
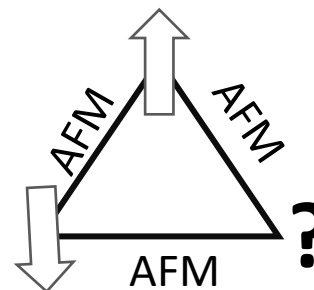
## **Initiation to the refinement of a magnetic Structures**

I will use two examples of Strongly frustrated Low-D inorganic Materials

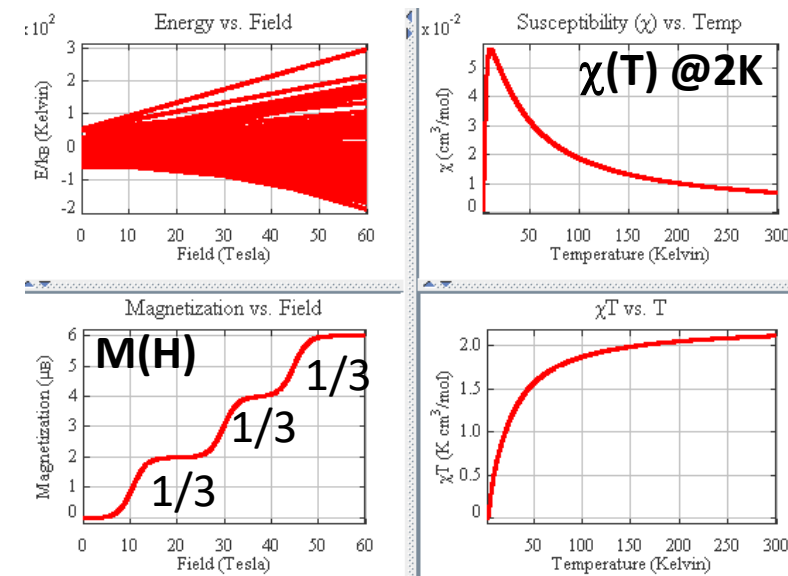
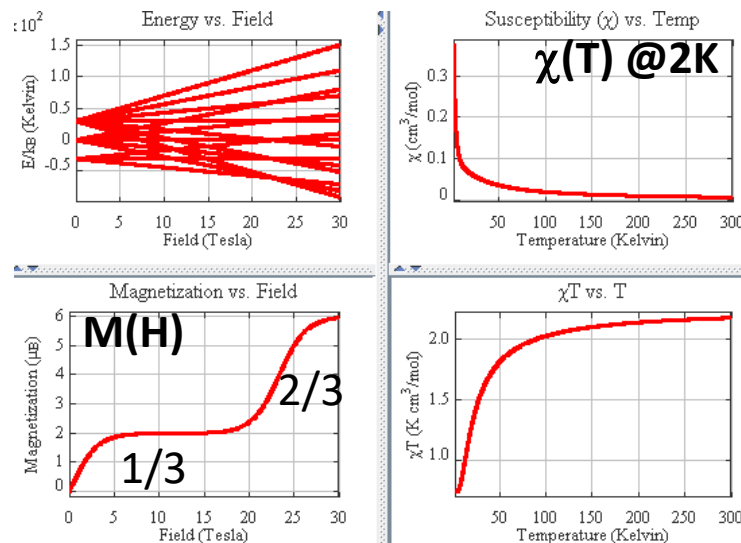
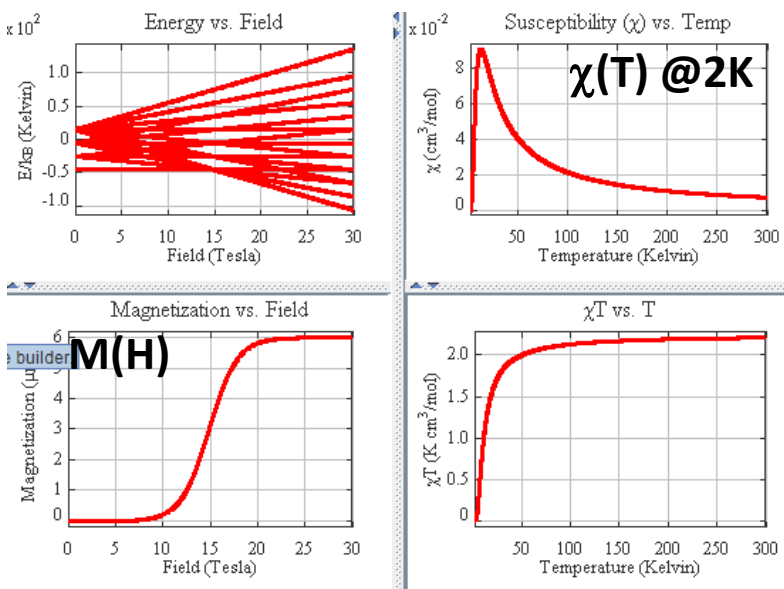
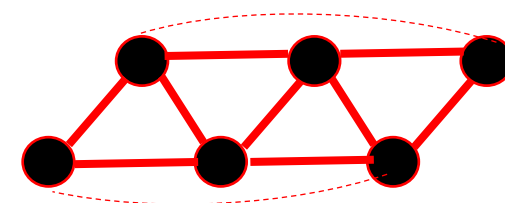
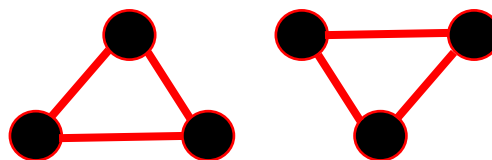
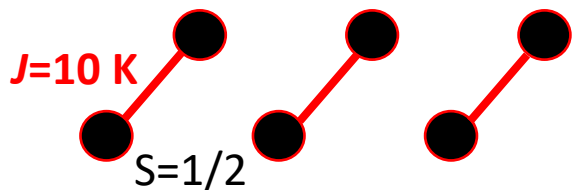
- 1) Frustration → original spin structures
- 2) Frustration → kills the magnetic ordering ... QSL state

# Frustration : the “standard model”

In most “real” materials it depends on the relative  $J_1, J_2, J_3$



Exact diagonalization,  $S=1/2$ , Heisenberg



Frustration effects :

1. “Fragilize” the AFM ground state

2. “Fragmentation” of  $M(H)$

# Neutron-atom elastic interaction

$$A = \underbrace{b}_{\text{nuclear origin / diffraction}} + \underbrace{2\mathbf{BI} \cdot \mathbf{S}}_{\text{Neutron spin and nuclear spin} \rightarrow \text{incoherent}} + \underbrace{(\gamma r_0/2)f(\mathbf{K})\mathbf{M}_{\perp} \cdot \mathbf{S}}_{\text{neutron spin and magn. moment / diffraction}}$$

*same order of magnitude*

$$\text{Int}_{\text{diff.}} = \underbrace{\{\mathbf{F}_N(\mathbf{K})\}^2}_{\text{cryst. Struct.}} + \underbrace{\{|\mathbf{F}_{\perp M}(\mathbf{K})|\}^2}_{\text{Magn. Struct.}}$$

Phys. IV France 11 (2001)  
 EDP Sciences, Les Ulis

**Structures magnétiques et symétries cristallines**

J. Schweizer

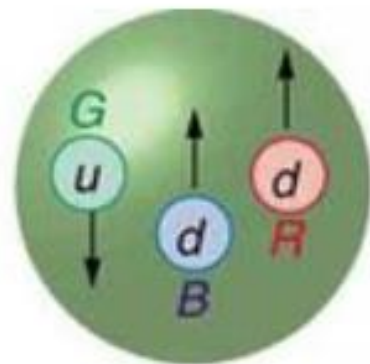
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Collection SFN 9 (2008) 65–85  
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 DOI: 10.1051/sfn:2008006

**Structures magnétiques, diffraction de neutrons et symétrie**

F. Bourée<sup>1</sup> et J. Rodri...

<sup>1</sup> LLB [CEA-CNRS], CEA/  
<sup>2</sup> Institut Laue-Langevin, E



Neutron 3 quarks

$$-\frac{1}{2} + \frac{1}{2} + \frac{1}{2} = \frac{1}{2}$$

*Crystallography on magnetic moments: Magn structure Factor :*

$$\mathbf{F}_{\perp M}(\mathbf{K}) = 0.27 \cdot 10^{-12} \sum_i f_j(\mathbf{K}) \{ \mathbf{m}_j(\mathbf{k}) - (\mathbf{e} \cdot \mathbf{m}_j(\mathbf{k})) \cdot \mathbf{e} \} \exp\{2\pi i \mathbf{K} \cdot \mathbf{r}_j\}$$

# Neutron-atom elastic interaction

$$A = \underbrace{b}_{\substack{\text{nuclear origin} \\ \text{/diffraction}}} + \underbrace{2\mathbf{BI} \cdot \mathbf{S}}_{\substack{\text{Neutron spin and nuclear spin} \\ \rightarrow \text{incoherent}}} + \underbrace{(\gamma r_0/2)f(\mathbf{K})\mathbf{M}_\perp \cdot \mathbf{S}}_{\substack{\text{neutron spin and magn.} \\ \text{moment/ diffraction}}}$$

*same order of magnitude*

$$\text{Int}_{\text{diff.}} = \underbrace{\{\mathbf{F}_N(\mathbf{K})\}^2}_{\text{cryst. Struct.}} + \underbrace{\{|\mathbf{F}_{\perp M}(\mathbf{K})|\}^2}_{\text{Magn. Struct.}}$$

## Magnetic Structure

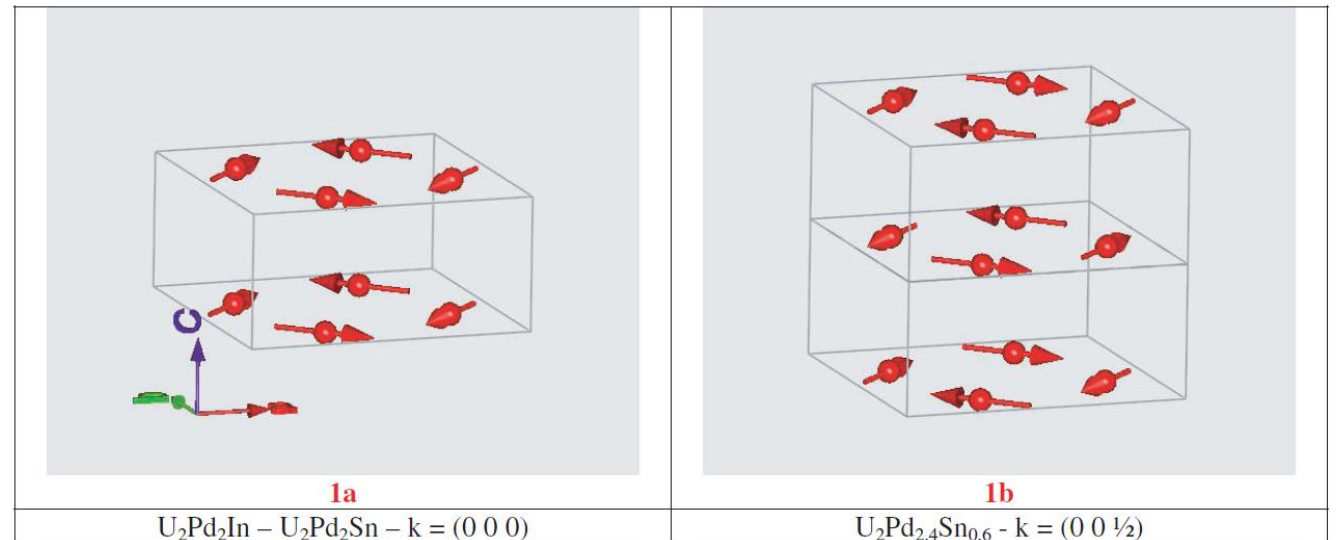
Below  $T_{\text{order}}$  the periodic magnetic struct. established will minimized the free energy of the system (m, J, k)

$\mathbf{K}$  = propagation vector / periodicity of the magn. structure

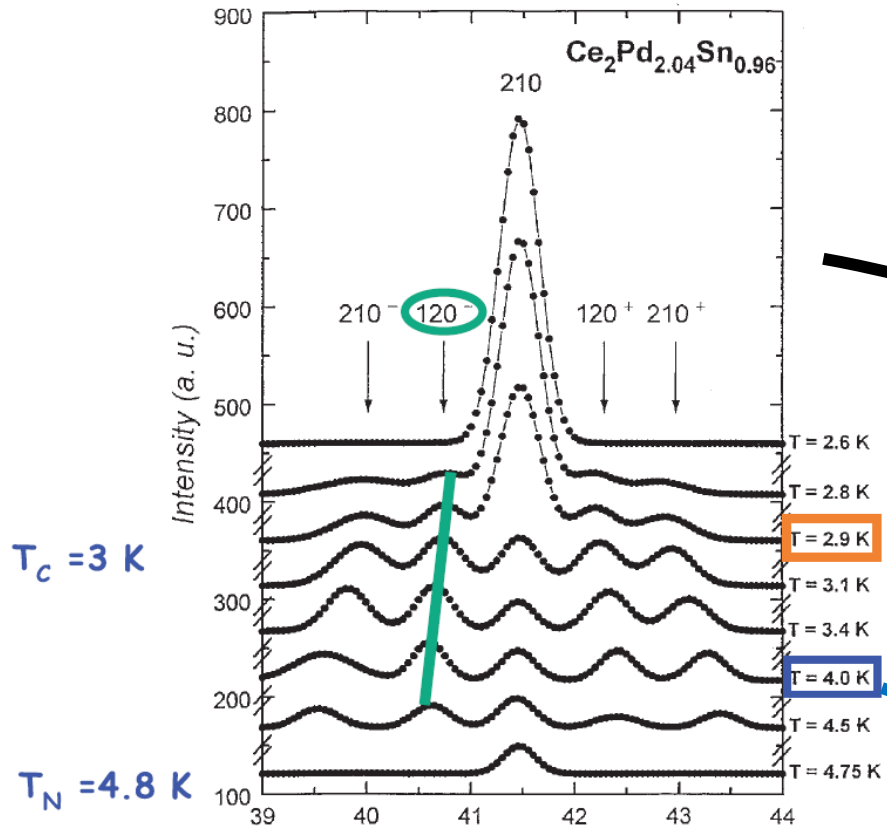
$\vec{\mathbf{K}}$  propagation vector

Supercell  $R^*$  vector, rational or incommensurate

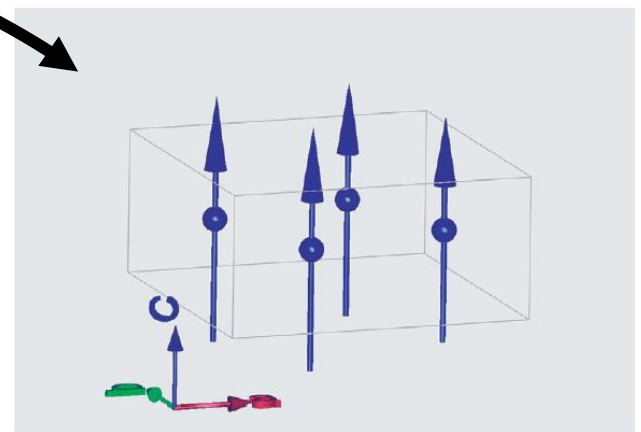
$$U = -N \sum_{j,j'} \sum_{\alpha,\beta} J_{jj'\alpha\beta}(\vec{\mathbf{k}}) m_{j\alpha}^{\vec{\mathbf{k}}} m_{j'\beta}^{-\vec{\mathbf{k}}}$$



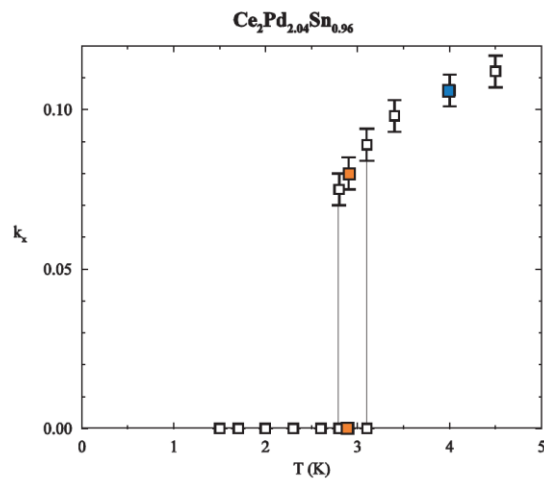
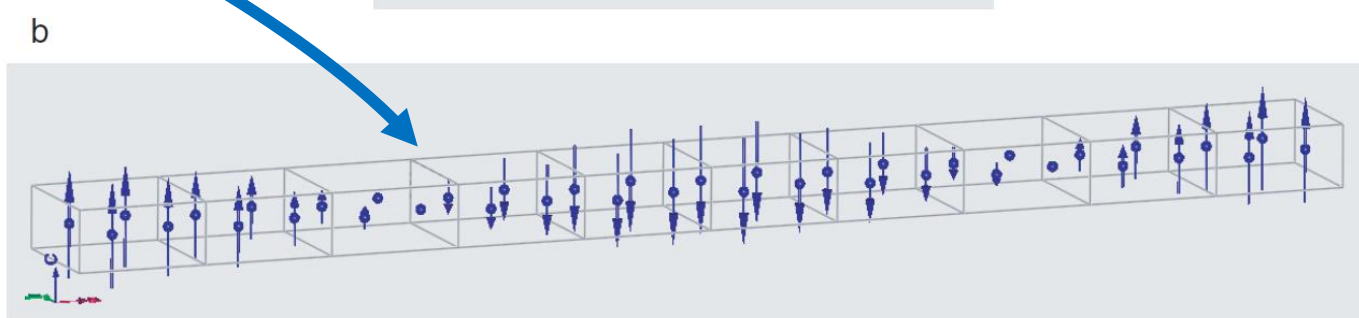
**Incommensurate magnetic structure** →  
spin density wave or cycloidal structure, helicoidal,  
conical, etc ...



$T = 2.6$  K



$T = 4$  K





# Solving a magnetic Structure → group Theory analysis (softwares : Sarah , Basireps ISOTROPY Software Suite, Bilbao Cryst....)

1•  $\vec{k}$  vector + Space group (G) → subgroup of k :  $G_k$

2•  $G_k$  + magn. atoms → magn. Matrix dim. 12

Ex :  $U_2Pd_2In$ ,  $k=(0,0,0)$  , G: P4/mbm =  $G_k$

4 atoms × 3 (Mx, My, Mz)

$U_1$	$x_U$	$\frac{1}{2} + x_U$	$\frac{1}{2}$
$U_2$	$1 - x_U$	$\frac{1}{2} - x_U$	$\frac{1}{2}$
$U_3$	$\frac{1}{2} - x_U$	$x_U$	$\frac{1}{2}$
$U_4$	$\frac{1}{2} + x_U$	$1 - x_U$	$\frac{1}{2}$

Notation	Loi de transformation (vecteur polaire)	Symbole (symétrie ponctuelle)	Translation	Situation	Notation Kovalev [13]
SYM( 1)	x, y, z	1			$h_1$
SYM( 2)	-x, -y, z	$2_z$			$h_4$
SYM( 3)	$\frac{1}{2}-x, \frac{1}{2}+y, -z$	$2_y$	$[0 \frac{1}{2} 0]$	$[\frac{1}{4} 0 0]$	$h_3$
SYM( 4)	$\frac{1}{2}+x, \frac{1}{2}-y, -z$	$2_x$	$[\frac{1}{2} 0 0]$	$[0 \frac{1}{4} 0]$	$h_2$
SYM( 5)	$\frac{1}{2}+y, \frac{1}{2}+x, -z$	$2_{[110]}$	$[\frac{1}{2} \frac{1}{2} 0]$		$h_{16}$
SYM( 6)	$\frac{1}{2}-y, \frac{1}{2}-x, -z$	$2_{[\bar{1}\bar{1}0]}$		$[0 \frac{1}{2} 0]$	$h_{13}$
SYM( 7)	y, -x, z	$4_z^3$			$h_{15}$
SYM( 8)	-y, x, z	$4_z$			$h_{14}$
SYM( 9)	-x, -y, -z	$\bar{1}$			$h_{25}$
SYM(10)	x, y, -z	$m_z$			$h_{28}$
SYM(11)	$\frac{1}{2}+x, \frac{1}{2}-y, z$	$m_y$	$[\frac{1}{2} 0 0]$	$[0 \frac{1}{4} 0]$	$h_{27}$
SYM(12)	$\frac{1}{2}-x, \frac{1}{2}+y, z$	$m_x$	$[0 \frac{1}{2} 0]$	$[\frac{1}{4} 0 0]$	$h_{26}$
SYM(13)	$\frac{1}{2}-y, \frac{1}{2}-x, z$	$m_{[110]}$		$[0 \frac{1}{2} 0]$	$h_{40}$
SYM(14)	$\frac{1}{2}+y, \frac{1}{2}+x, z$	$m_{[\bar{1}\bar{1}0]}$	$[\frac{1}{2} \frac{1}{2} 0]$		$h_{37}$
SYM(15)	-y, x, -z	$4_z^3$			$h_{39}$
SYM(16)	y, -x, -z	$4_z$			$h_{38}$

		U1			U2			U3			U4		
		Mx	My	Mz	Mx	My	Mz	Mx	My	Mz	Mx	My	Mz
U1	Mx	-1											
	My		-1										
	Mz			1									
U2	Mx				-1								
	My					-1							
	Mz						1						
U3	Mx							-1					
	My								-1				
	Mz									1			
U4	Mx										-1		
	My											-1	
	Mz												1



z

3 • 16 symmetries  $\rightarrow$  16 magn. Matrix of dim. 12

- All together  $\rightarrow$  representation  $\Gamma$  of the  $G_k$  group of dim. 12
- The magn repr.  $\Gamma$  can be reduced in irreducible representation  $\Gamma_{IR}$  of various dimensions

*Symmetry*

P4/mbm	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16		
$\Gamma_1$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1		
$\Gamma_2$	1	1	1	1	1	1	1	1	-1	-1	$D_{3h}$	$E$	$2C_3$	$3C_2$	$\sigma_h$	$2S_3$	$3\sigma_v$	
$\Gamma_3$	1	1	1	1	-1	-1	-1	-1	1	1	$A'_1$	1	1	1	1	1	1	$x^2 + y^2, z^2$
$\Gamma_4$	1	1	1	1	-1	-1	-1	-1	-1	-1	$A'_2$	1	1	-1	1	1	-1	$R_z$
$\Gamma_5$	1	1	-1	-1	1	1	-1	-1	1	1	$E'$	2	-1	0	2	-1	0	$(x, y)$
$\Gamma_6$	1	1	-1	-1	1	1	-1	-1	-1	-1	$A''_1$	1	1	1	-1	-1	-1	$(x^2 - y^2, xy)$
$\Gamma_7$	1	1	-1	-1	-1	-1	1	1	1	1	$A''_2$	1	1	-1	-1	-1	1	$z$
$\Gamma_8$	1	1	-1	-1	-1	-1	1	1	-1	-1	$E''$	2	-1	0	-2	1	0	$(R_x, R_y)$
$\Gamma_9$	1 0 0 1	-1 0 0 -1	1 0 0 -1	-1 0 0 1	0 1 1 0	0 -1 -1 0	0 -1 1 0	0 1 -1 0	-1 0 0 -1	1 0 0 1	0 1 0 -1	0 -1 0 1	-1 0 1 0	0 -1 -1 0	0 -1 1 0	0 1 -1 0		
$\Gamma_{10}$	1 0 0 1	-1 0 0 -1	1 0 0 -1	-1 0 0 1	0 1 1 0	0 -1 -1 0	0 -1 1 0	0 1 -1 0	1 0 0 1	-1 0 0 -1	1 0 0 -1	-1 0 0 1	0 1 1 0	0 -1 -1 0	0 -1 1 0	0 1 -1 0		

- by orthogonality it can be decomposed in  $n \Gamma_{IR}$ , each associated with an order  $m \longrightarrow \Gamma = \Gamma_2 \oplus \Gamma_3 \oplus \Gamma_4 \oplus \Gamma_6 \oplus \Gamma_7 \oplus \Gamma_8 \oplus \Gamma_9 \oplus 2\Gamma_{10}$ .

$$\Gamma = \Gamma_2 \oplus \Gamma_3 \oplus \Gamma_4 \oplus \Gamma_6 \oplus \Gamma_7 \oplus \Gamma_8 \oplus \Gamma_9 \oplus 2\Gamma_{10}$$

Each of the  $\Gamma_{IR}$  gives a magnetic structure and needs to be tested against the diffr. data

For instance  $\Gamma_2$

**U1** SYMM x,y,z

0.2700 0.7700 0.5000

sk(1): (u,u,0)

**U2** SYMM -x,-y,z

0.2700 -0.7700 0.5000

sk(2): (-u,-u,0)

**U3** SYMM -x+1/2,y+1/2,-z

0.2300 1.2700 -0.5000

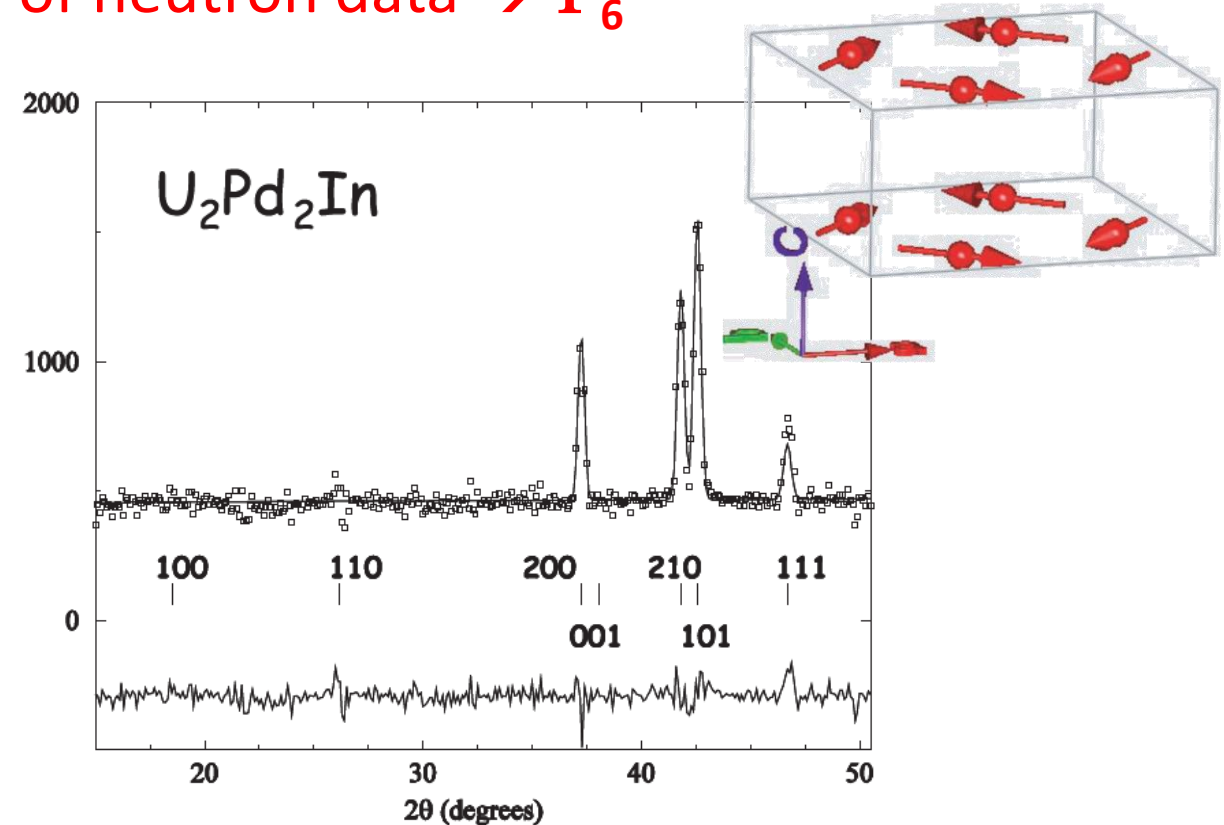
sk(3): (-u,u,0)

**U4** SYMM x+1/2,-y+1/2,-z

0.7700 -0.2700 -0.5000

sk(4): (u,-u,0)

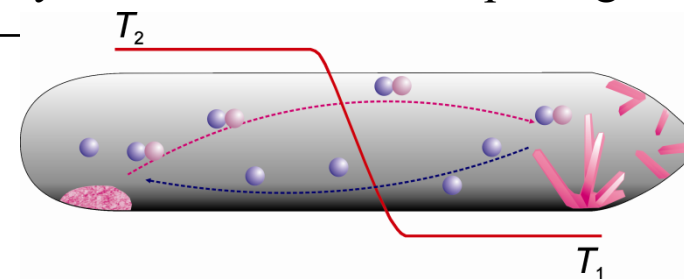
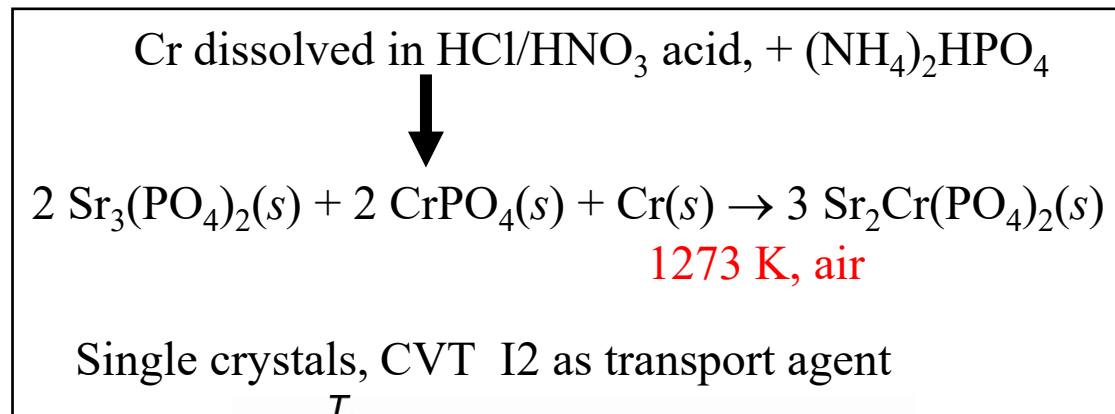
Rietveld refinement  
of neutron data  $\rightarrow \Gamma_6$



# Frustration drives idle spin in $\text{Sr}_2\text{Cr}(\text{PO}_4)_2$

$\text{Cr}^{2+}$ , d4, S=2, L=2, strong JT effect  
 Unstable in oxide, easily oxidized  
 CrO does not exist  $\rightarrow$   $\text{Cr}_2\text{O}_3$  too stable

$\text{Cr}^{2+}$  is stabilized by oxo-anions ( $\text{SiO}_4^{4-}$ ,  $\text{BO}_4^{5-}$ ,  $\text{PO}_4^{3-}$  ...  $\text{F}^-$ ) by inductive effect, e.g  $\text{CrF}_2$   
 $\text{Cr}^{2+} \gg \text{O} \gg \text{P}$



# $\text{Sr}_2\text{Cr}(\text{PO}_4)_2$ : Crystal structure

**P b c a (61)**

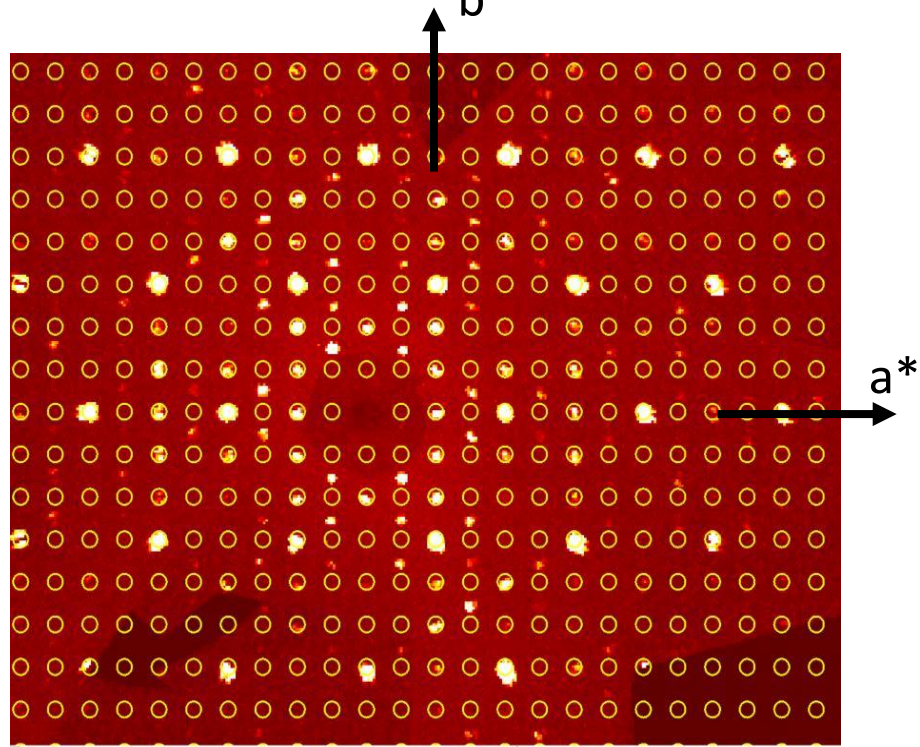
$a=10.7064(6) \text{ \AA}$

$b=9.2730(5) \text{ \AA}$

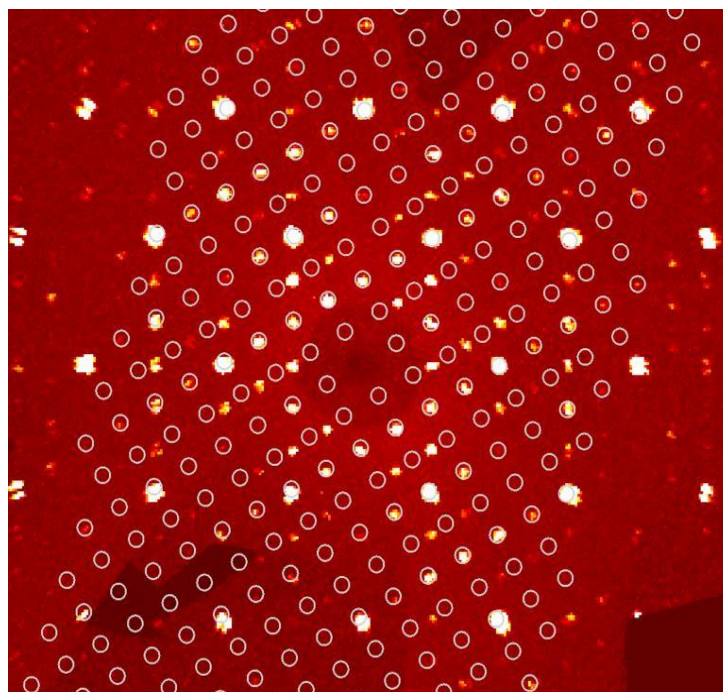
$c=21.2720(7)$

- Orthorhombic Pbc a
- Two individual layers
- Pseudo trigonal

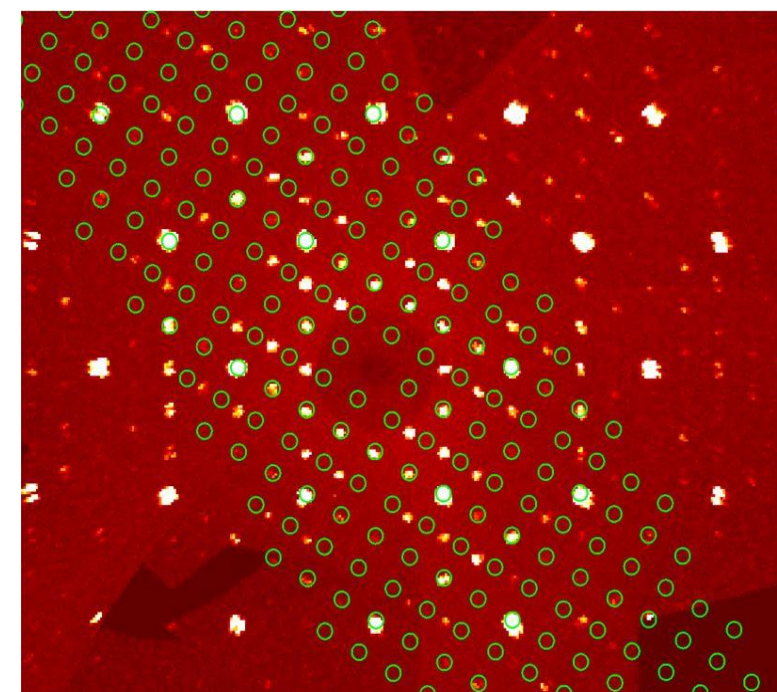
$hk0$  – domain 1



$hk0$  – domain 2

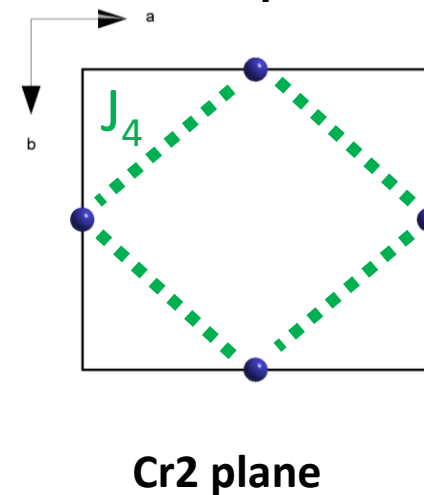
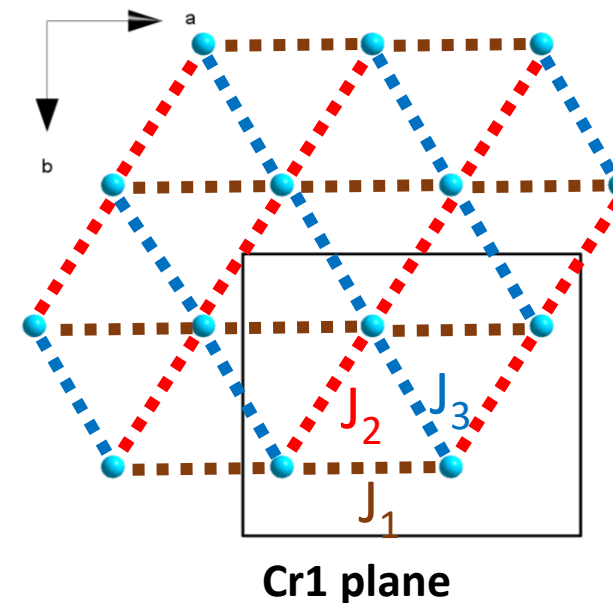
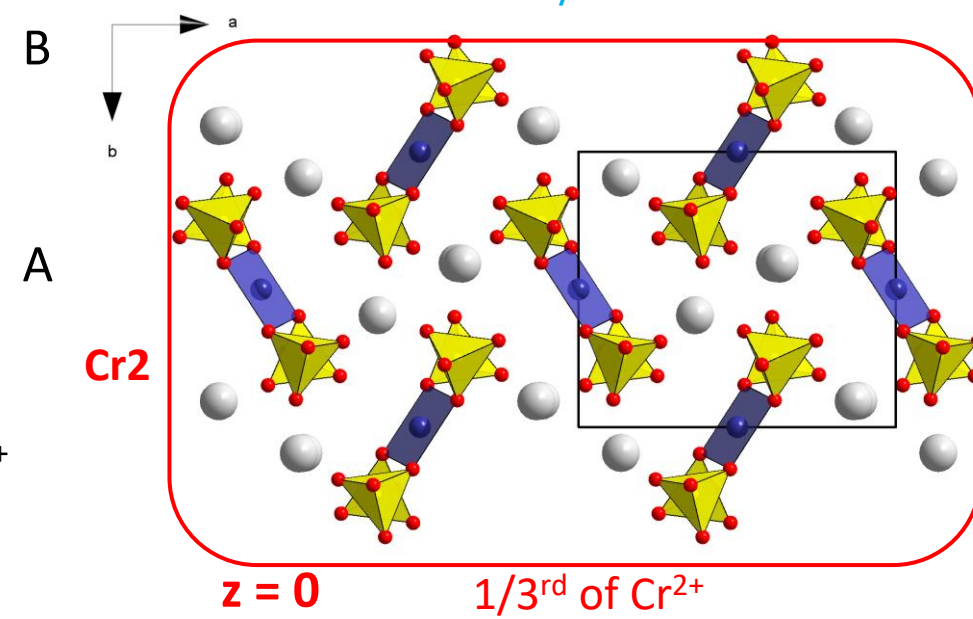
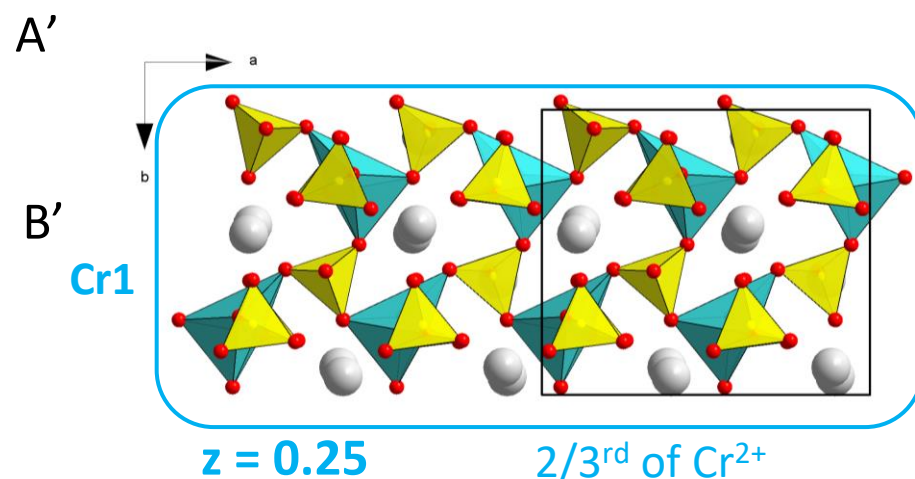
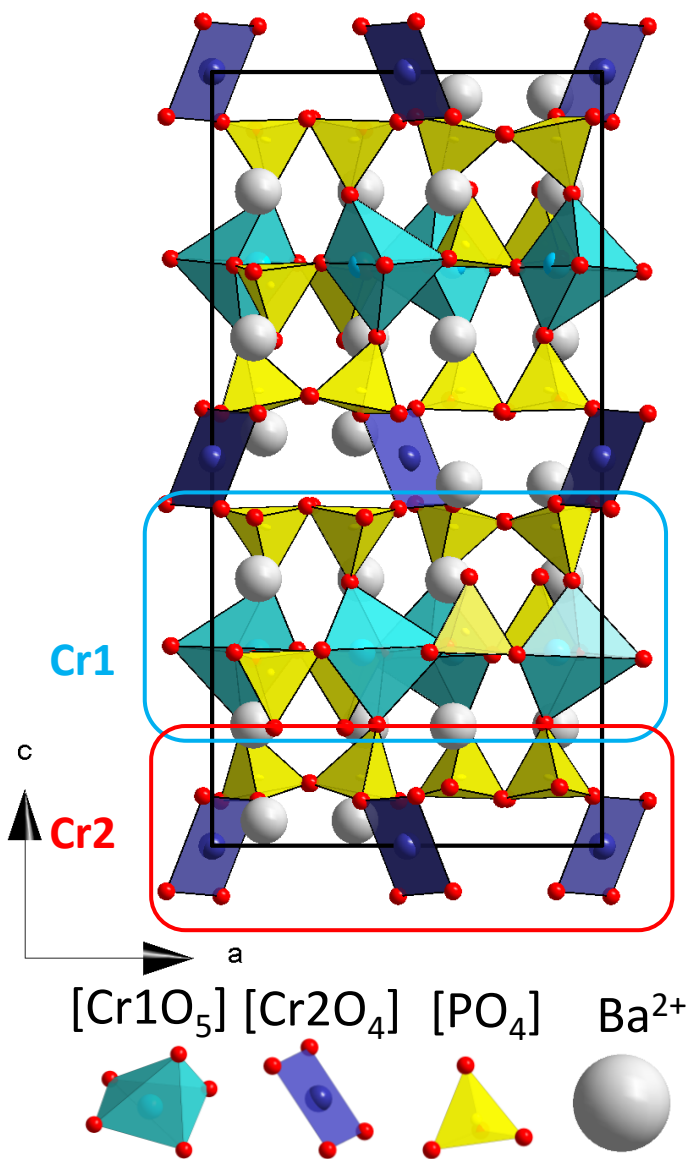


$hk0$  – domain 3



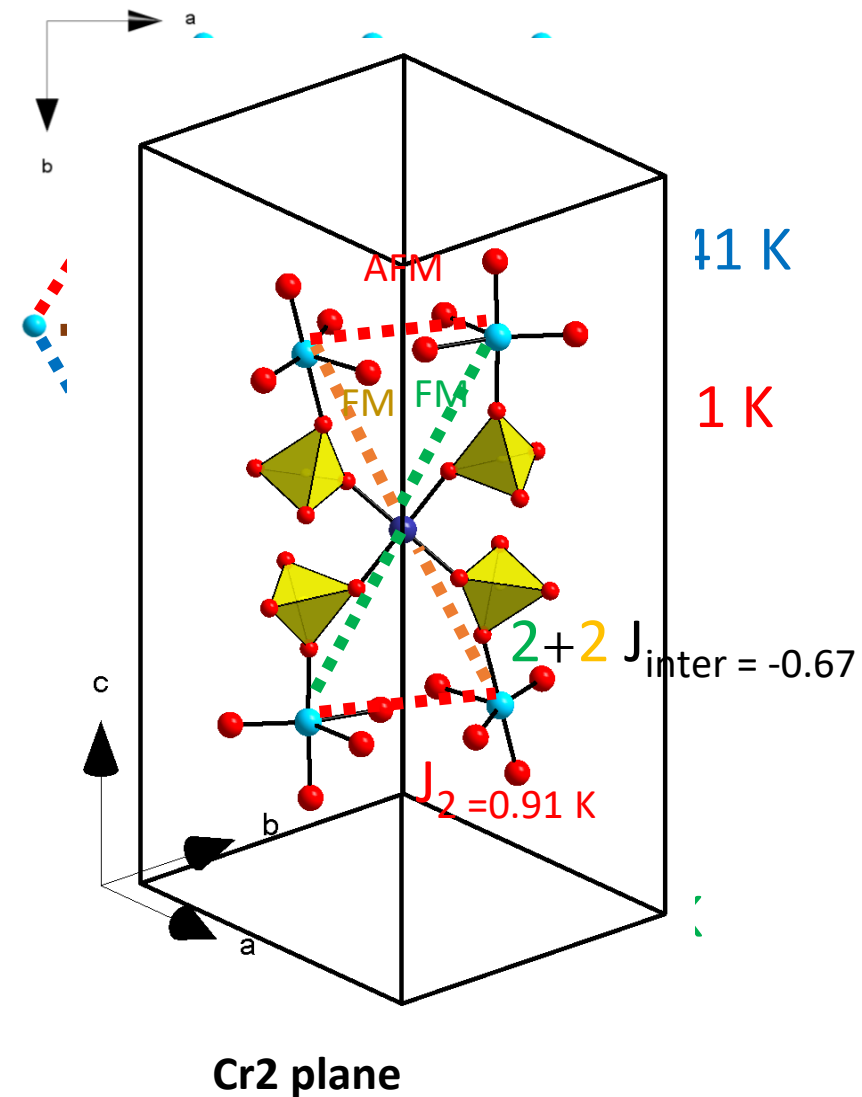
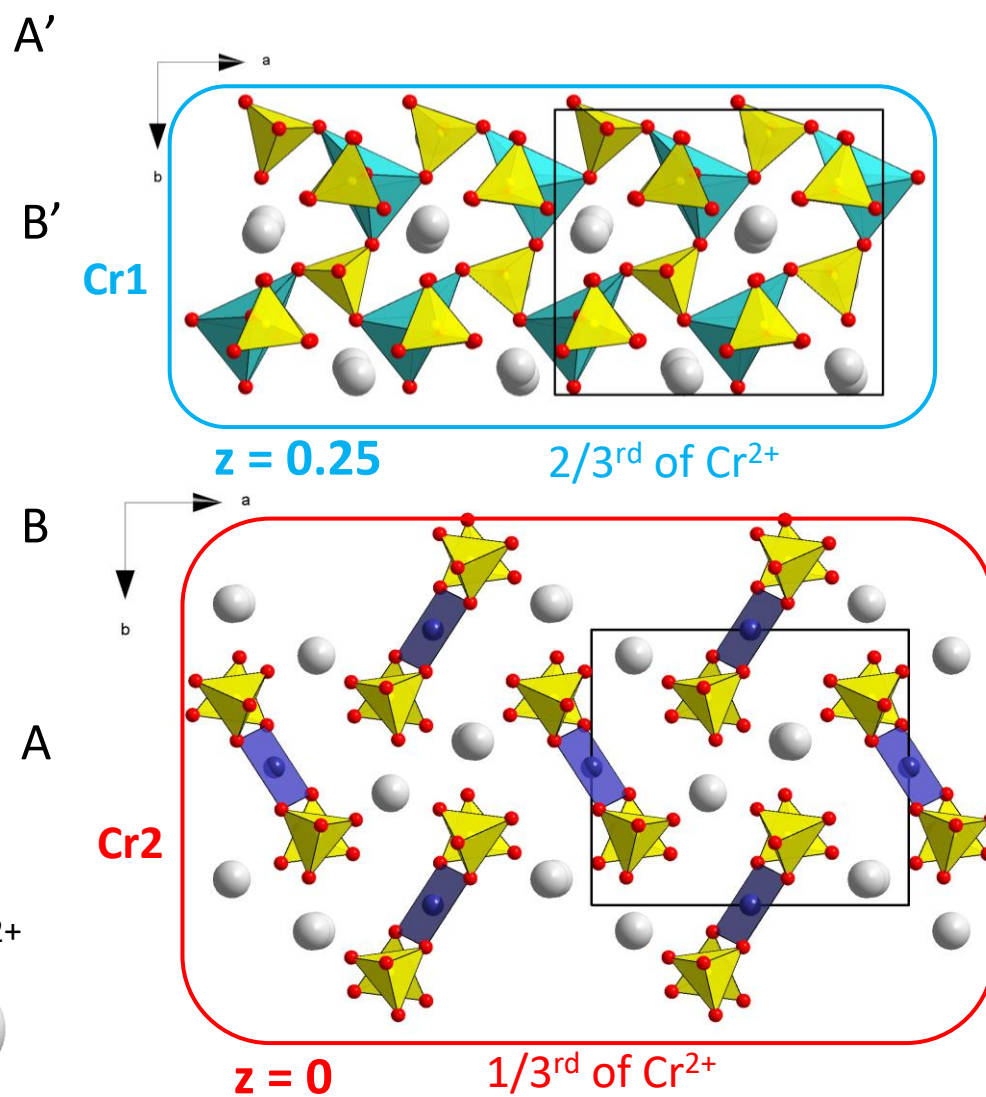
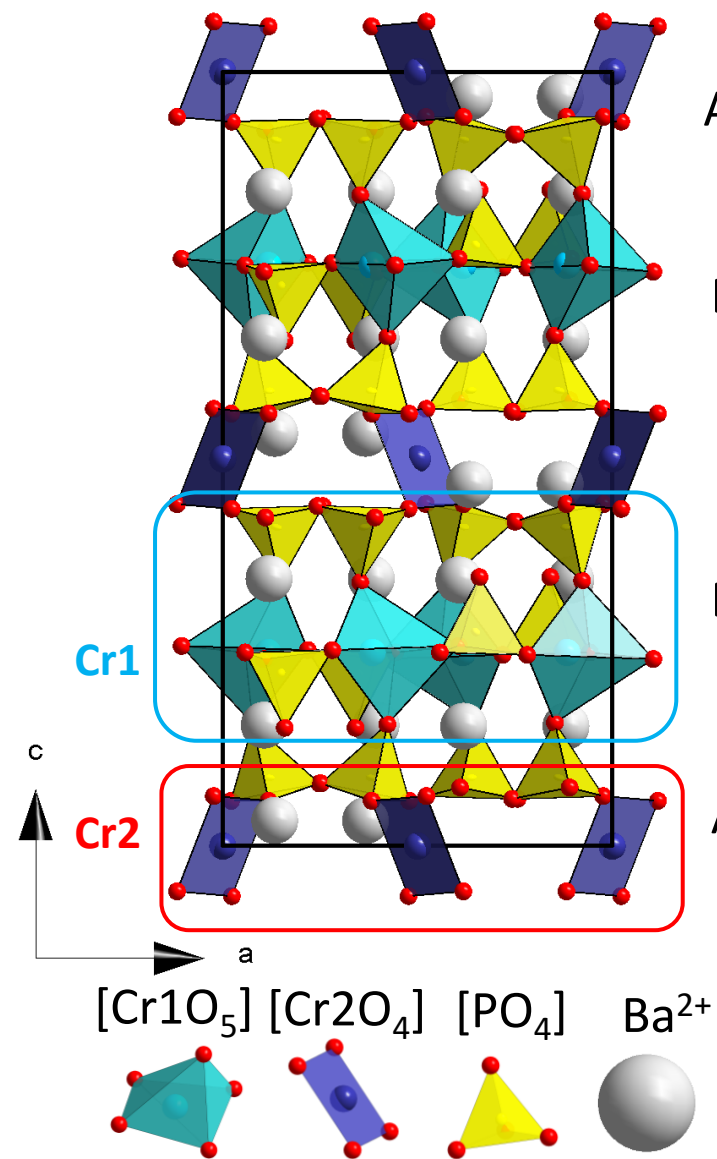
# Sr<sub>2</sub>Cr(PO<sub>4</sub>)<sub>2</sub> : Crystal structure

- Orthorhombic Pbc<sub>a</sub>
- Two individual layers
- Pseudo trigonal

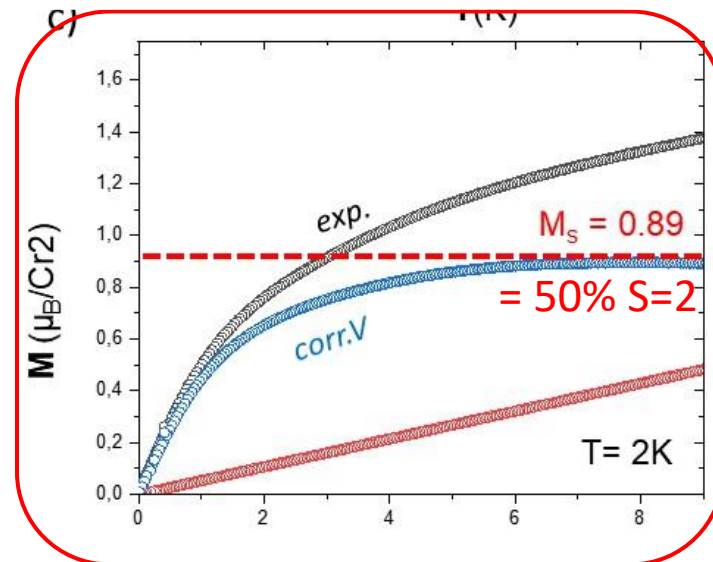
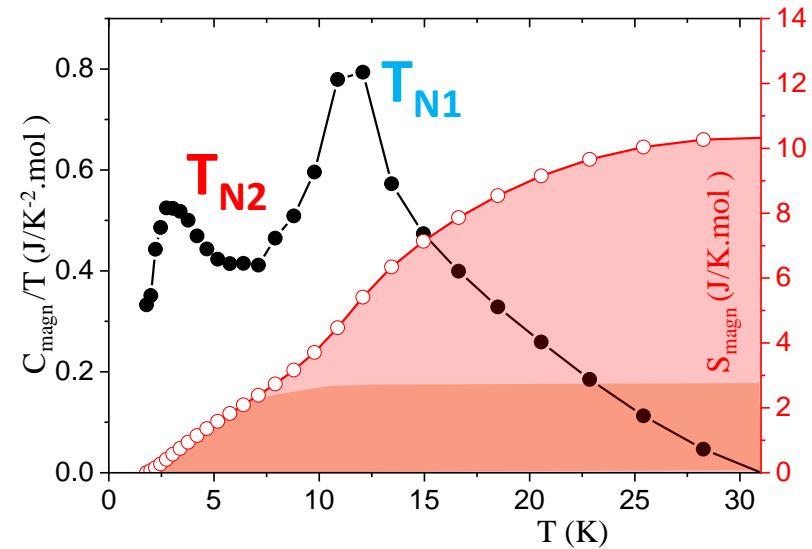
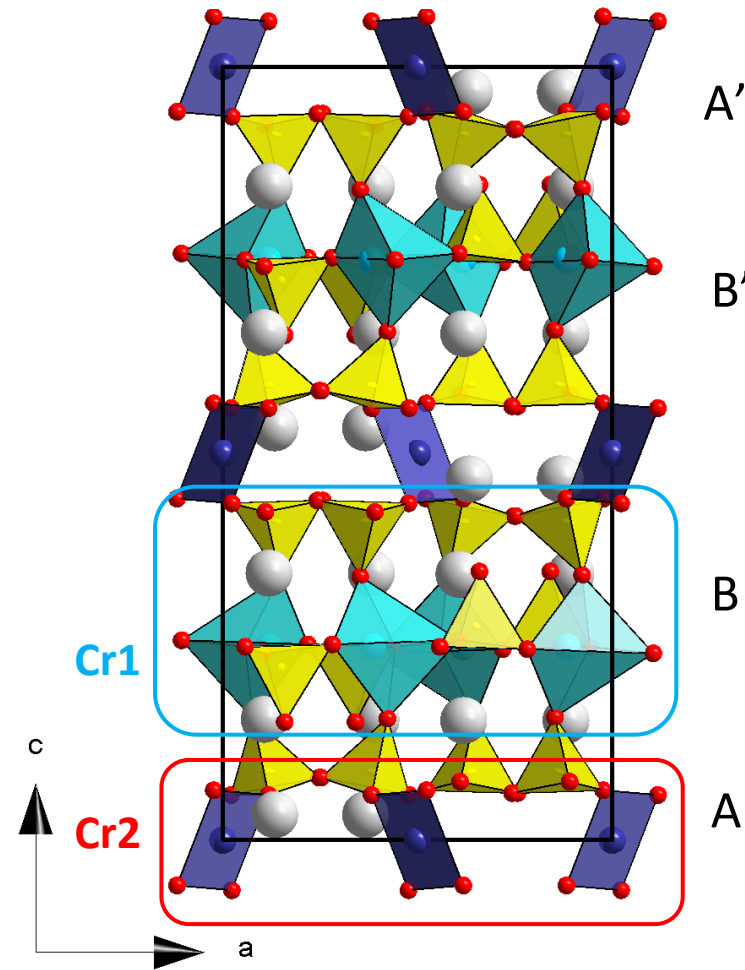


# Sr<sub>2</sub>Cr(PO<sub>4</sub>)<sub>2</sub> : Crystal structure

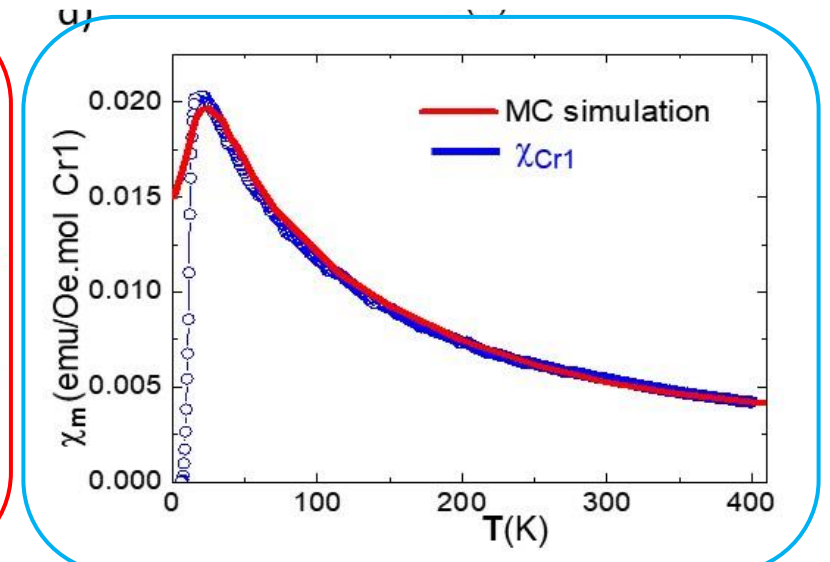
- Orthorhombic Pbc<sub>a</sub>
- Two individual layers
- Pseudo trigonal



# Sr<sub>2</sub>Cr(PO<sub>4</sub>)<sub>2</sub> : The contribution of the two sublattices



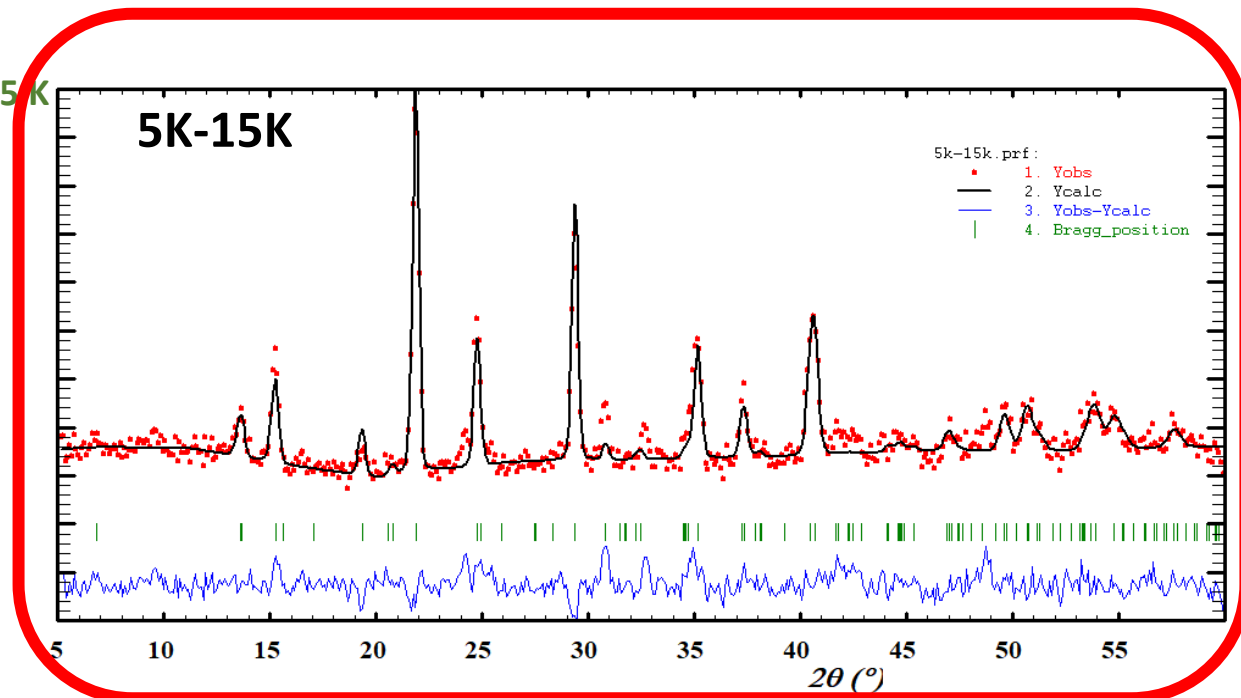
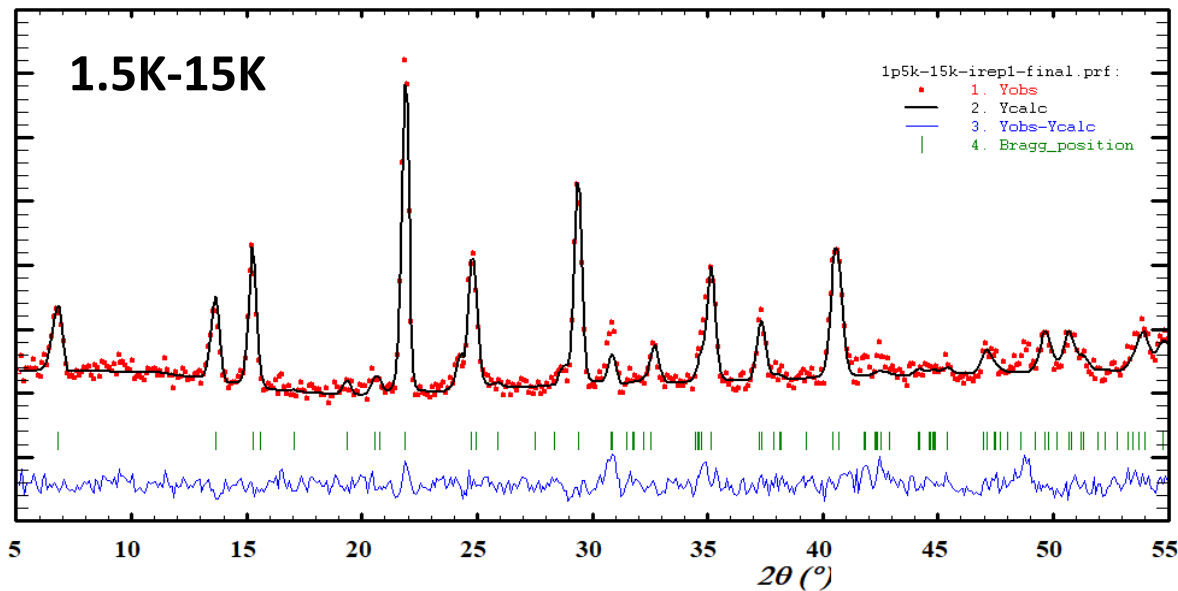
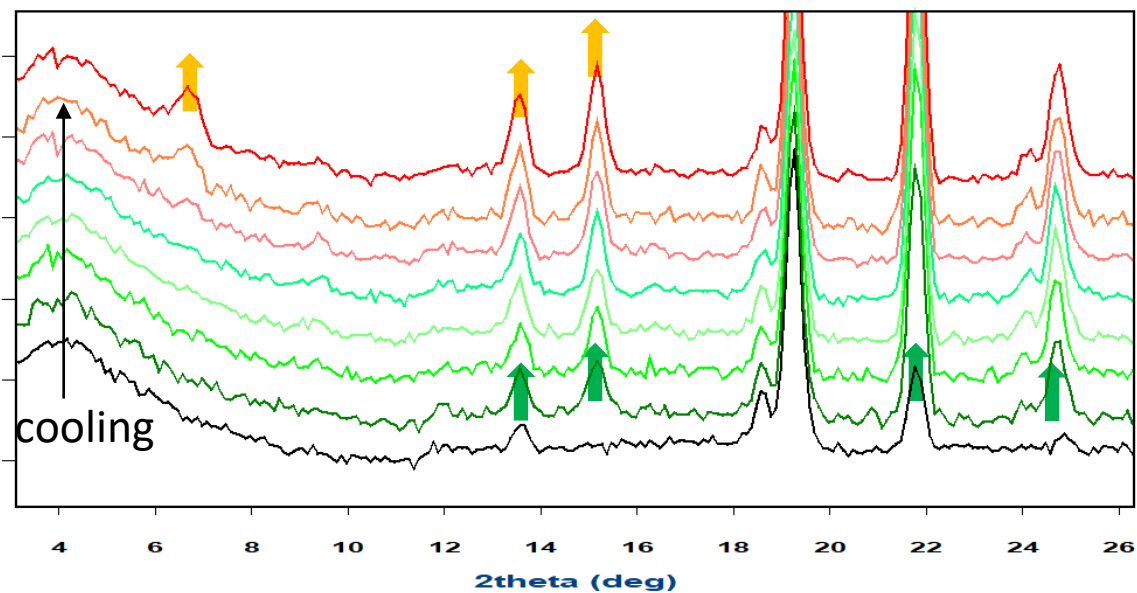
Low T: 50% only of the Cr2 seems fluctuating, 50% ordered



Cr1 AFM-2D lattice ordering at 13.6 K + Cr2 idle = "quasi-paramagn."



NPD, D1B ILL



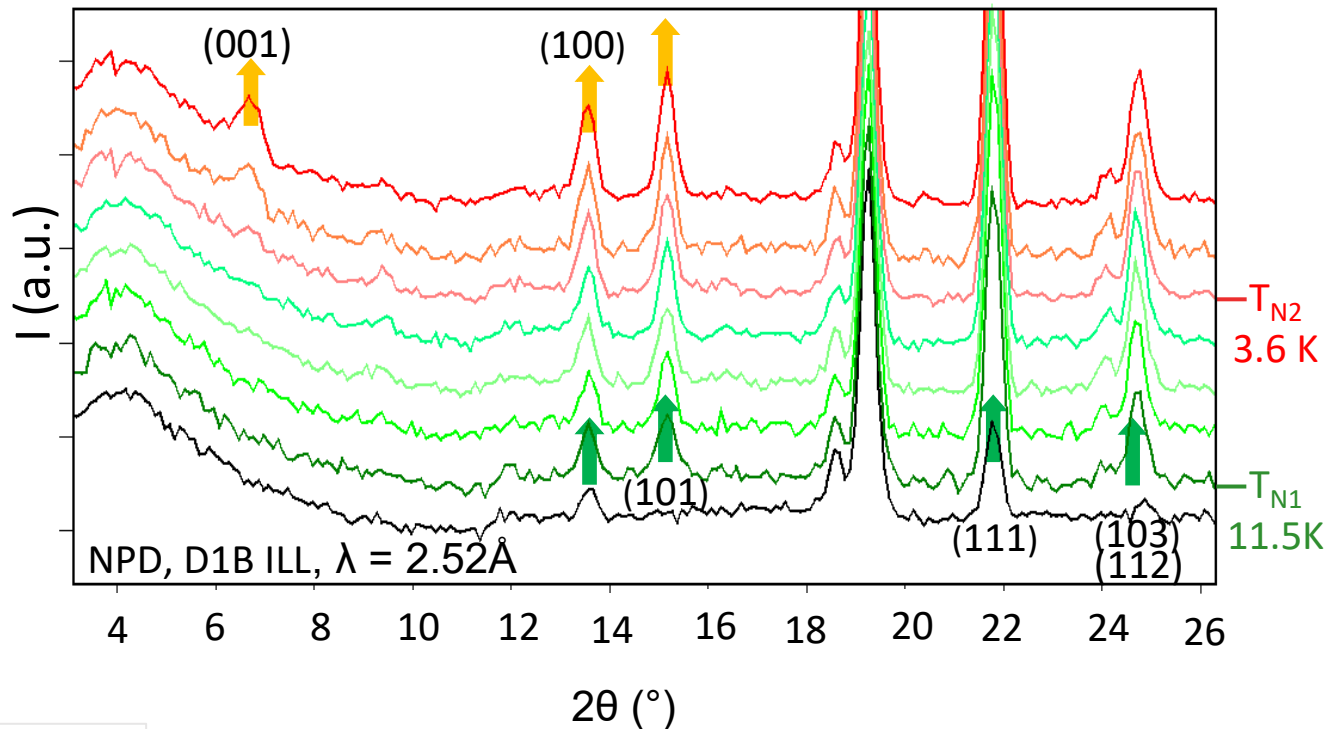
# Solving the Magnetic Structure at 5K

**1<sup>st</sup> step** : find the propagation k vector

All magn. peaks are indexable in the crystal cell

$$\rightarrow \mathbf{K} = (0,0,0)$$

**2<sup>nd</sup> step** : Group theory analyzis



```
=====
CALCULATIONS FOR SITE : Cr(1)
=====
```

=> Decomposition of the Magnetic/Mechanic representation:

-> GAMMA(Magnetic): 3 Irep\_k( 1) + 3 Irep\_k( 2) + 3 Irep\_k( 3) + 3 Irep\_k( 4) + 3 Irep\_k( 5) + 3 Irep\_k( 6) + 3 Irep\_k( 7) + 3 Irep\_k( 8)

```
=====
CALCULATIONS FOR SITE : Cr(2)
=====
```

=> Decomposition of the Magnetic/Mechanic representation:

-> GAMMA(Magnetic): 3 Irep\_k( 1) + 3 Irep\_k( 3) + 3 Irep\_k( 5) + 3 Irep\_k( 7)

*demo*

## What represent these IRreps ?

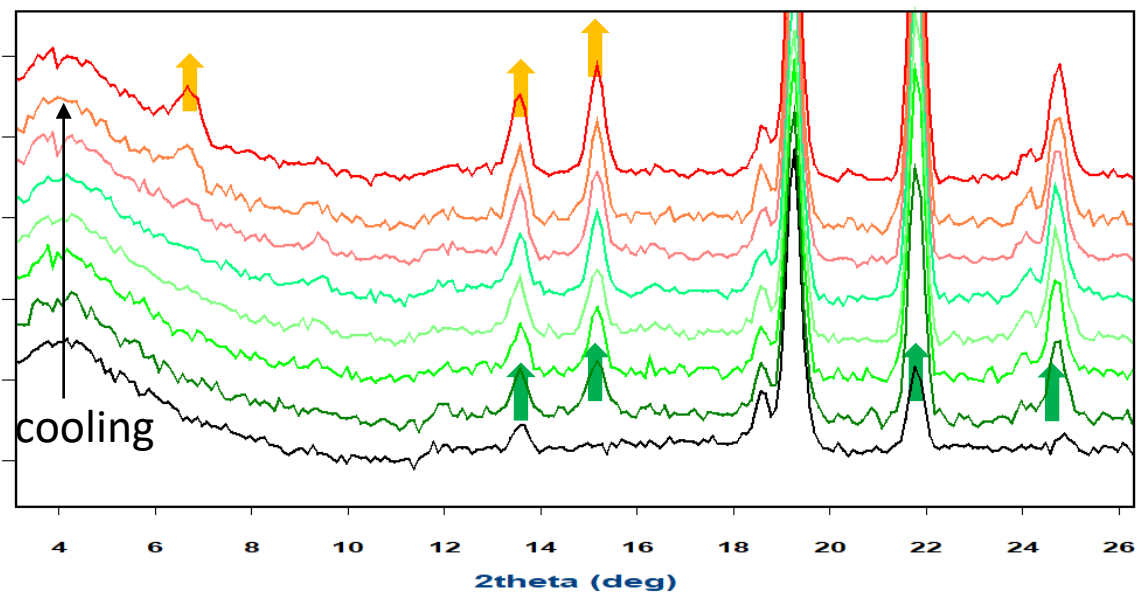
=> Basis functions of Representation IRrep( 8) of dimension 1 contained 3 times in GAMMA  
Representation number : 8 for Site: 1  
Number of basis functions: 3

```
----- Block-of-lines for PCR start just below this line
P -1                                <--Space group symbol for hkl generation
! Nsym   Cen   Laue Ireps N_Bas
   8     1     1    -1     3
! Real(0)-Imaginary(1) indicator for Ci
  0  0  0
SYMM x,y,z      Mx      My      Mz
BASR  1  0  0  0  1  0  0  0  1
BASR  0  0  0  0  0  0  0  0  0
SYMM -x+1/2,-y,z+1/2
BASR  1  0  0  0  1  0  0  0 -1
BASR  0  0  0  0  0  0  0  0  0
SYMM -x,y+1/2,-z+1/2
BASR  1  0  0  0 -1  0  0  0  1
BASR  0  0  0  0  0  0  0  0  0
SYMM x+1/2,-y+1/2,-z
BASR  1  0  0  0 -1  0  0  0 -1
BASR  0  0  0  0  0  0  0  0  0
SYMM -x,-y,-z
BASR -1  0  0  0 -1  0  0  0 -1
BASR  0  0  0  0  0  0  0  0  0
SYMM x+1/2,y,-z+1/2
BASR -1  0  0  0 -1  0  0  0  1
BASR  0  0  0  0  0  0  0  0  0
SYMM x,-y+1/2,z+1/2
BASR -1  0  0  0  1  0  0  0 -1
BASR  0  0  0  0  0  0  0  0  0
SYMM -x+1/2,y+1/2,z
BASR -1  0  0  0  1  0  0  0  1
BASR  0  0  0  0  0  0  0  0  0
----- End-of-block of lines for PCR
```

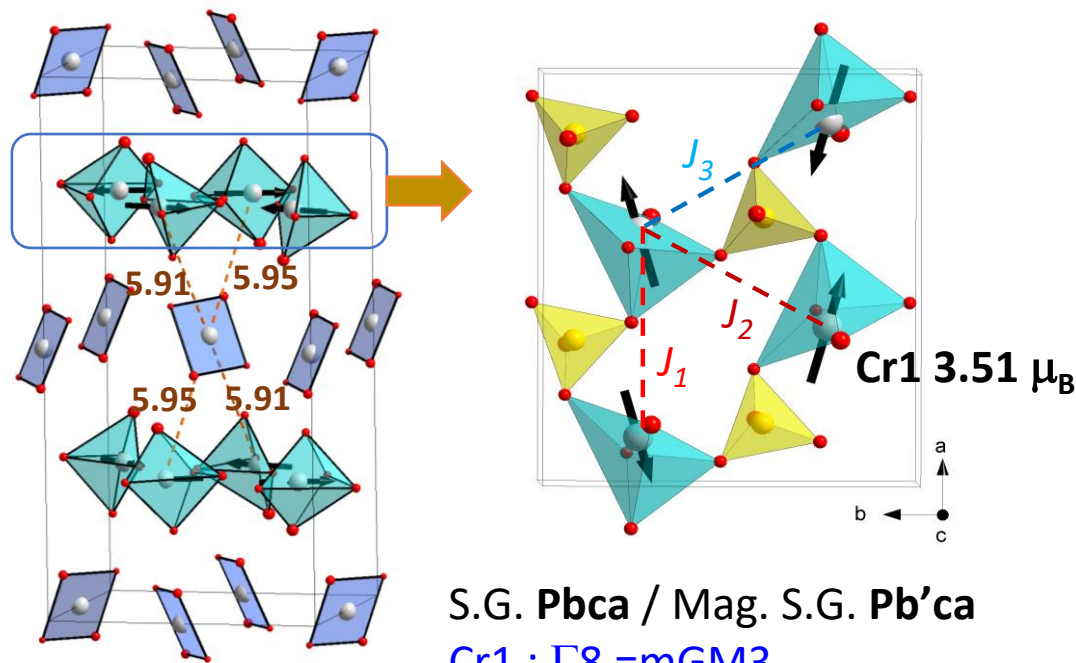
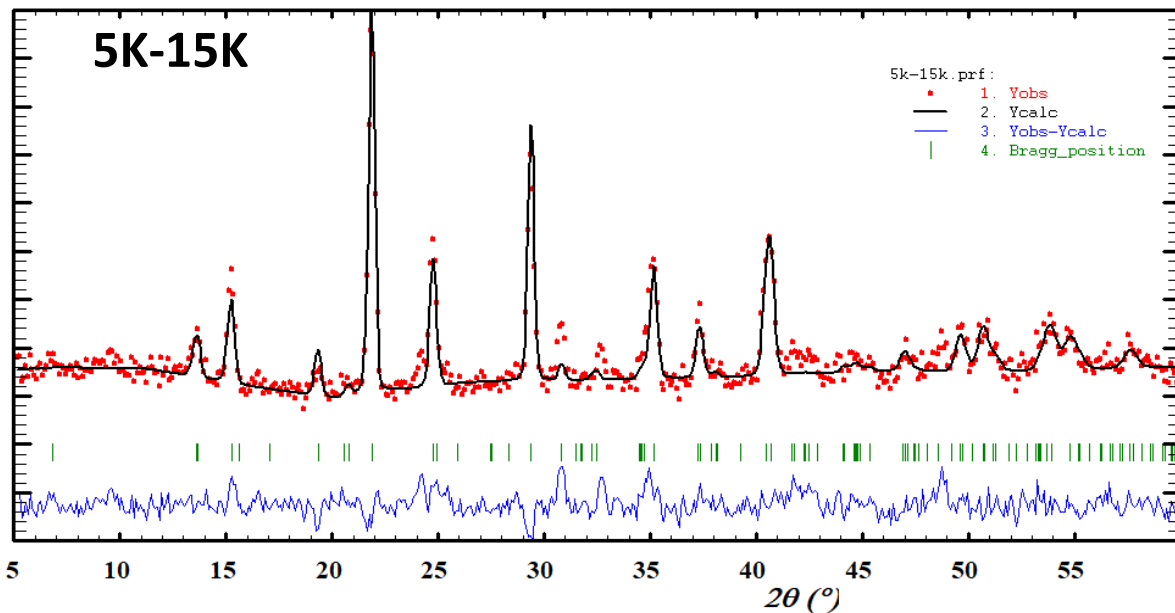
**3<sup>rd</sup> step** : test all IRreps → refinement of magnetic structure

*demo*

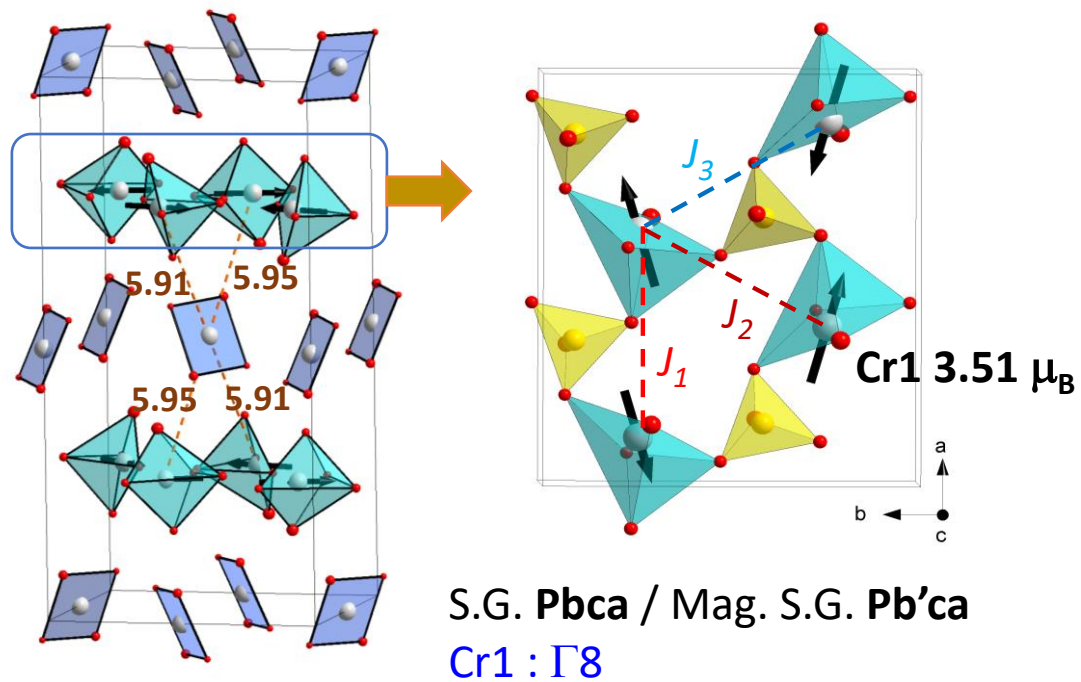
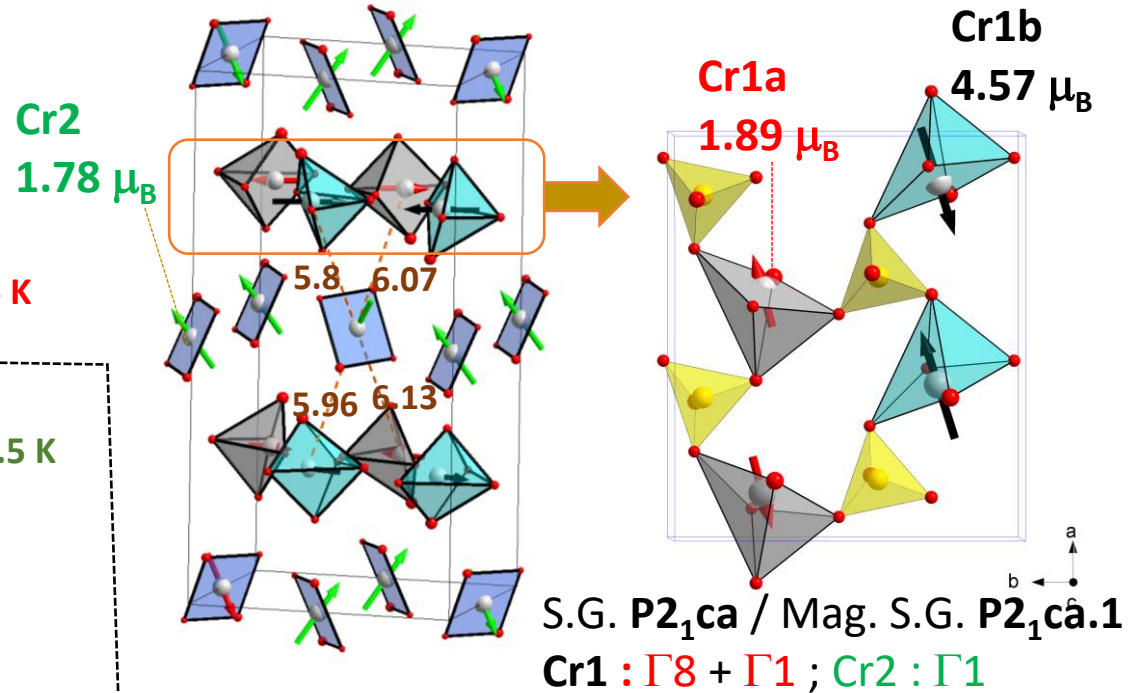
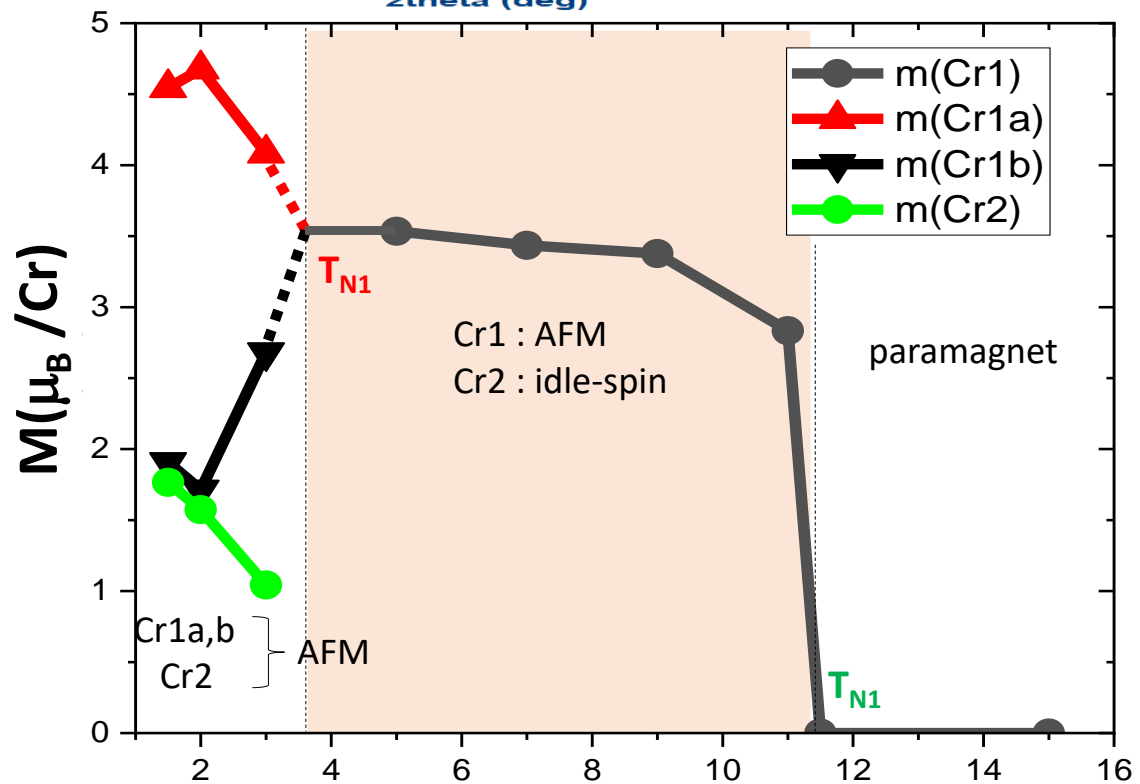
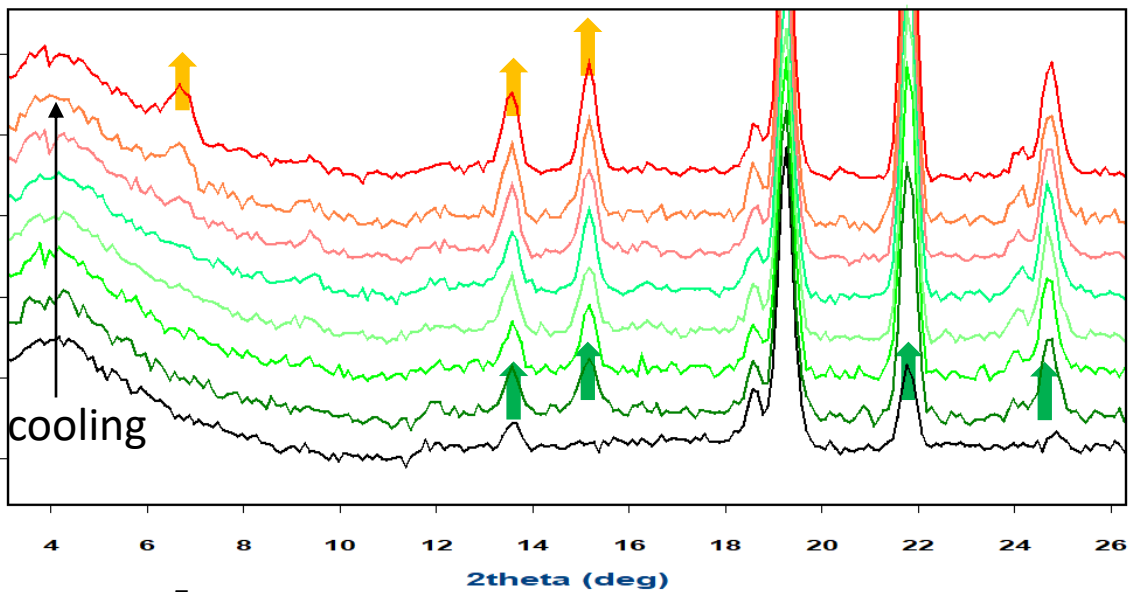
NPD, D1B ILL

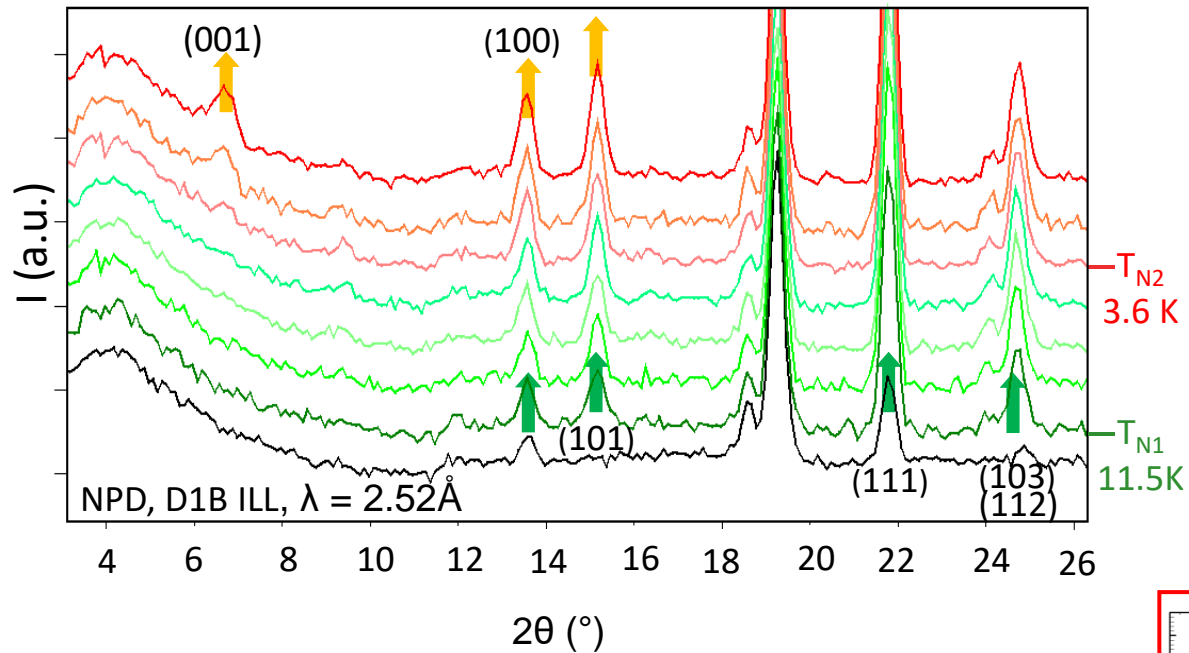


Cr(1) AFM-2D lattice ordering  
between 13.6 K - 3.6 K  
+ Cr2 idle = "quasi-paramagn."

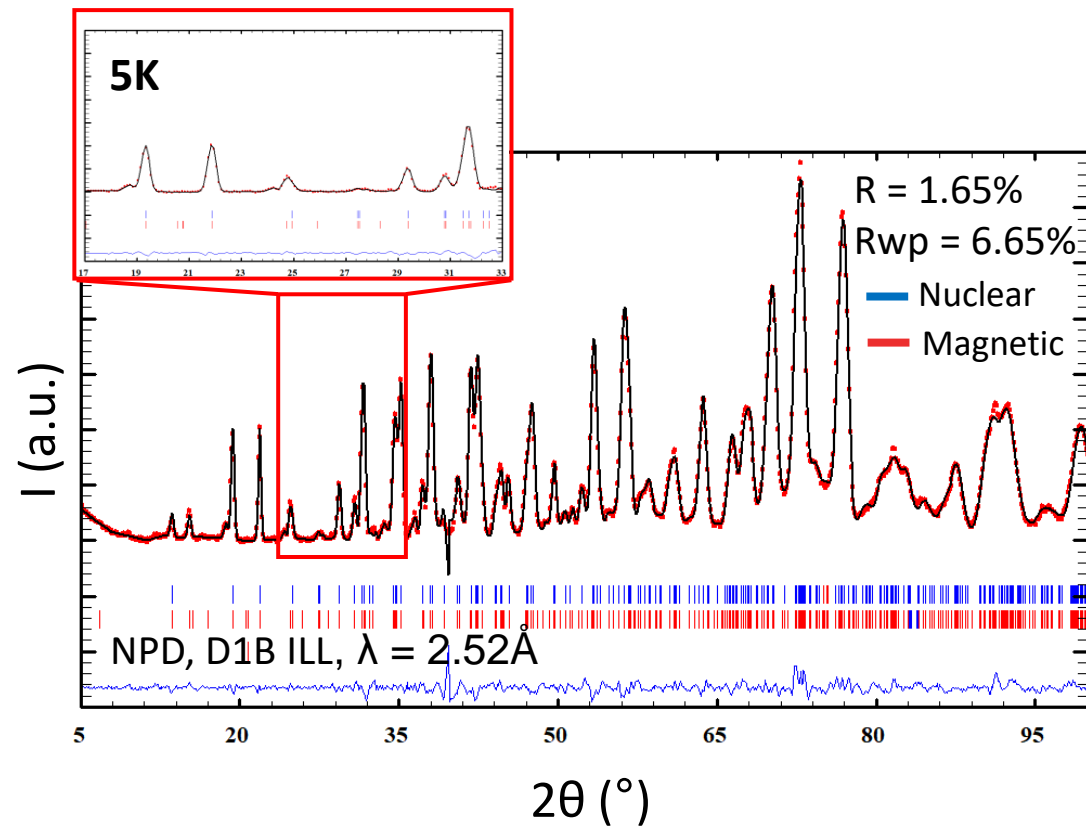
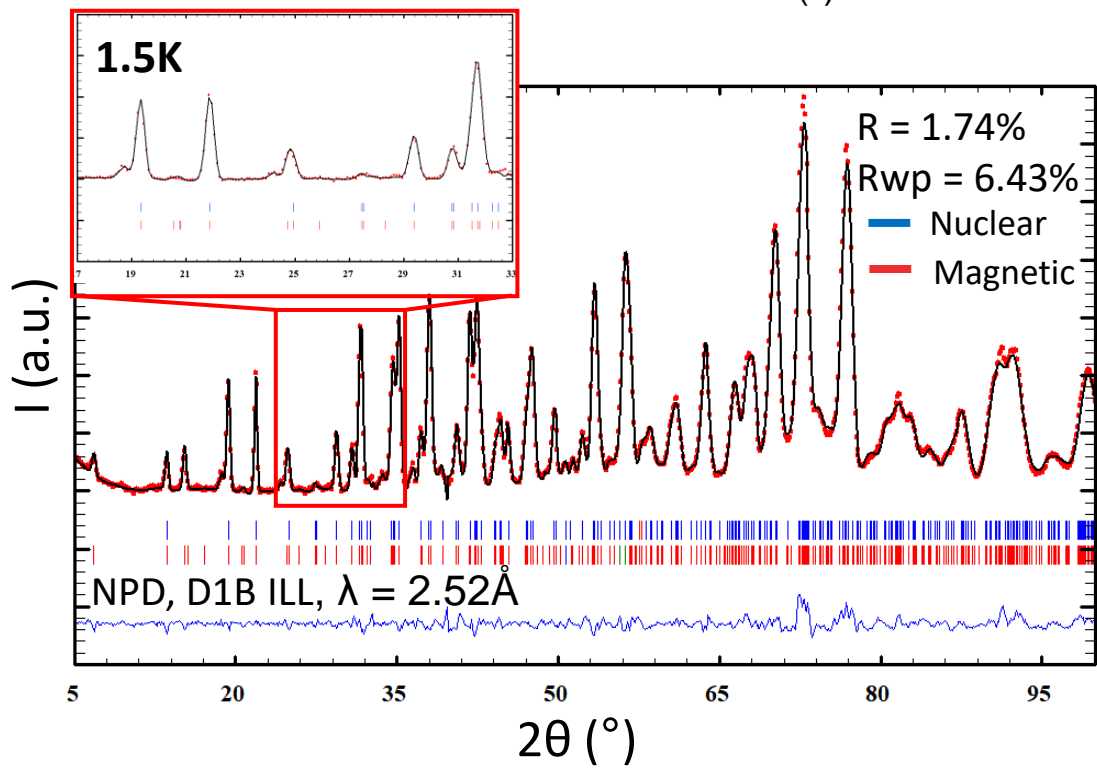


# NPD, D1B ILL





We can also refine the full NPD pattern : nuclear + magnetic



# Frustration ... Towards « true » Quantum Spin Liquid

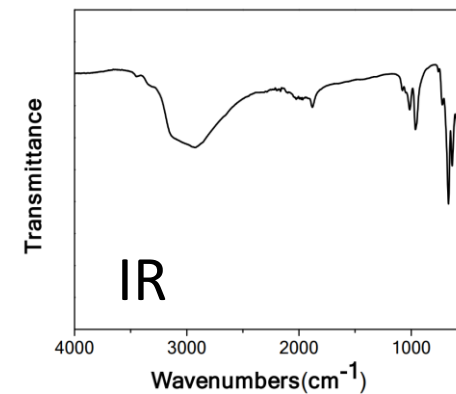
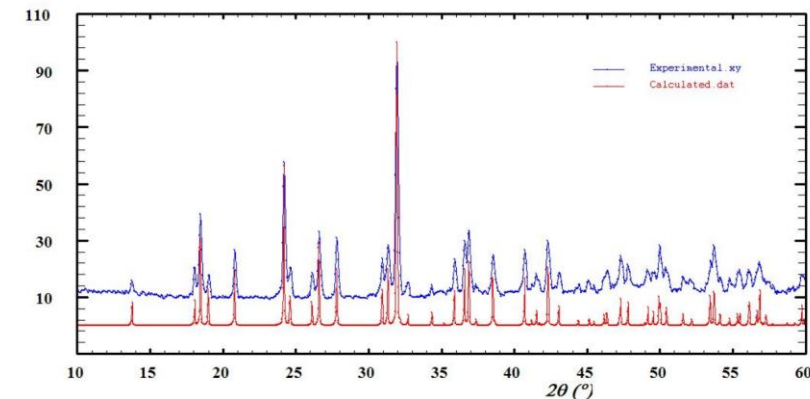
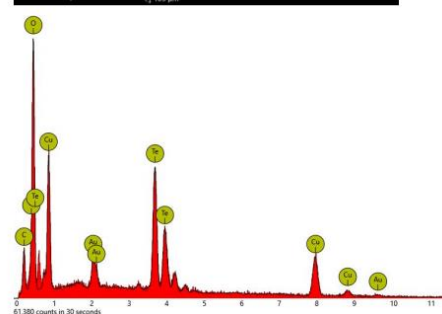
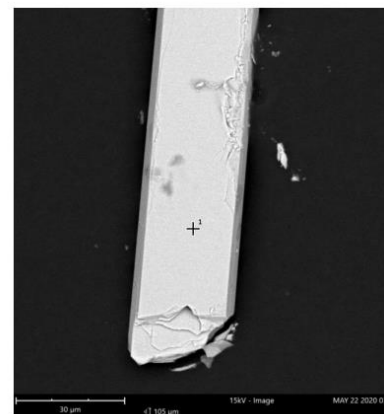
## Synthesis

### • *The $Cu_3Te_2O_5(OH)_4$ compound*



- Teflon-lined bomb : autogeneous pressure
- Hydrazine red. agent: stabilization low-redox
- Fluoride-rich media (HF)

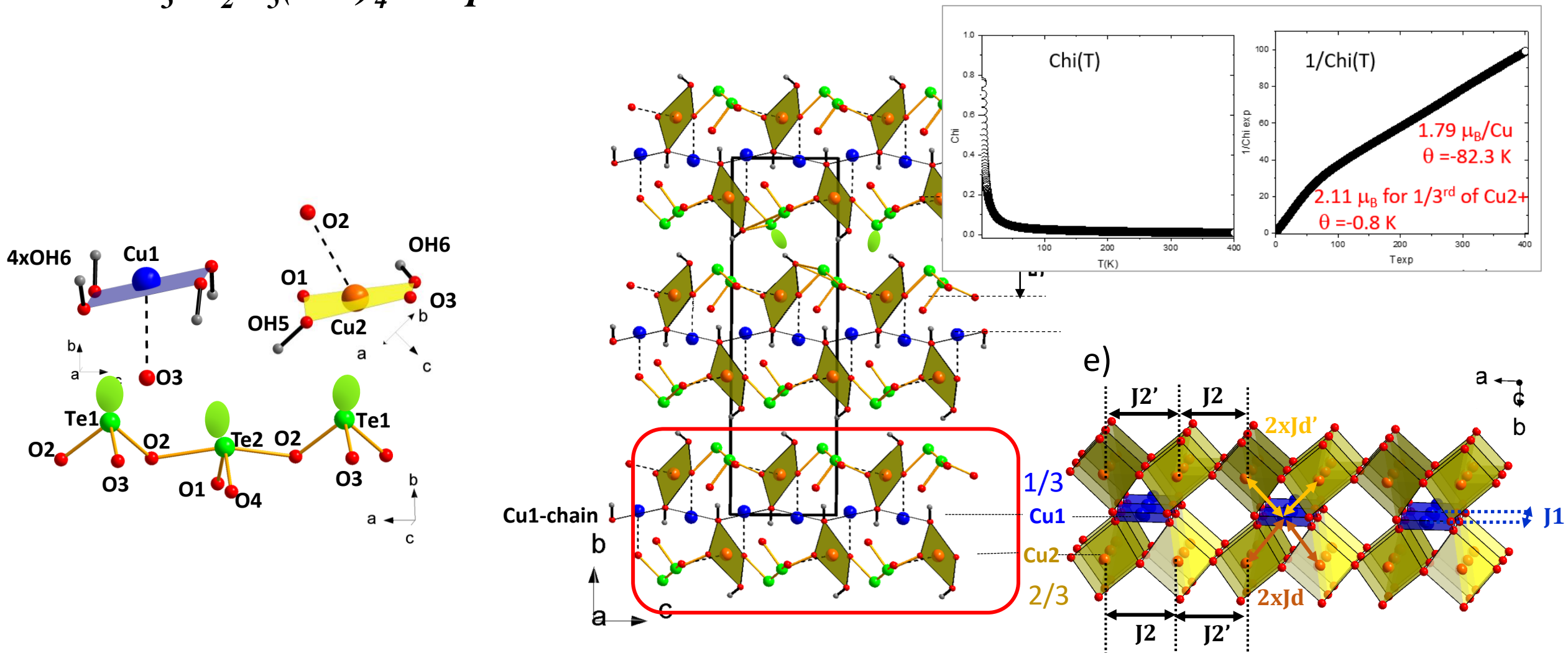
$Cu_2(OH)_2CO^3$  (Alfa Aesar, >55%),  $TeO_2$  (Alfa Aesar, 99%), HF (Energy Chemical, 48% wt in  $H_2O$ ), Hydrazine monohydrate, (Sinopharm, >85%)



# Frustration ... Towards « true » Quantum Spin Liquid

*... requires « perfect » frustrated Spins ...*

## • The $Cu_3Te_2O_5(OH)_4$ compound

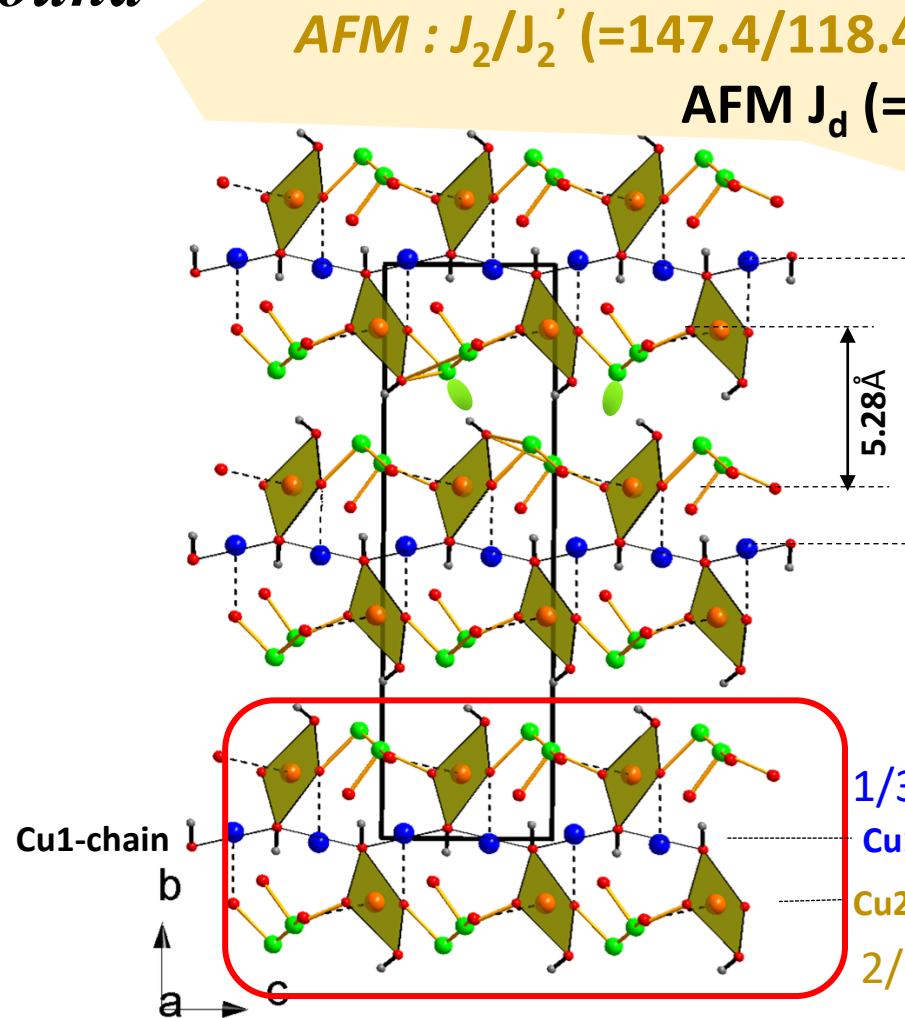
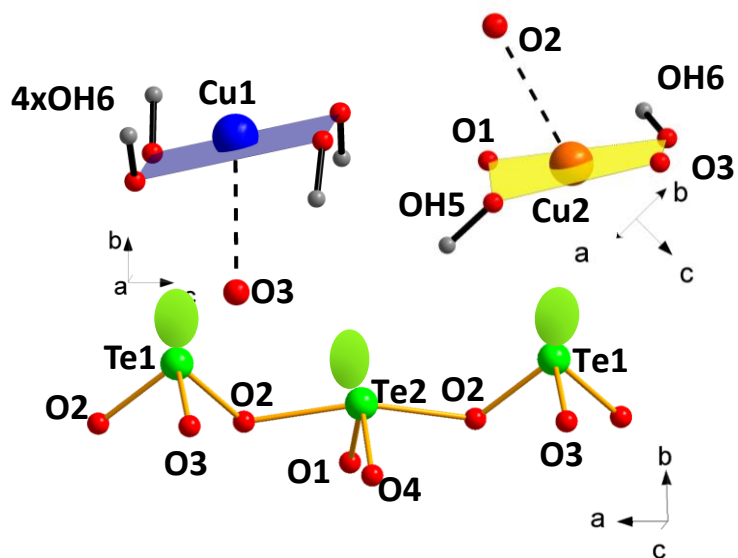




# Frustration ... Towards « true » Quantum Spin Liquid

*... requires « perfect » frustrated Spins ...*

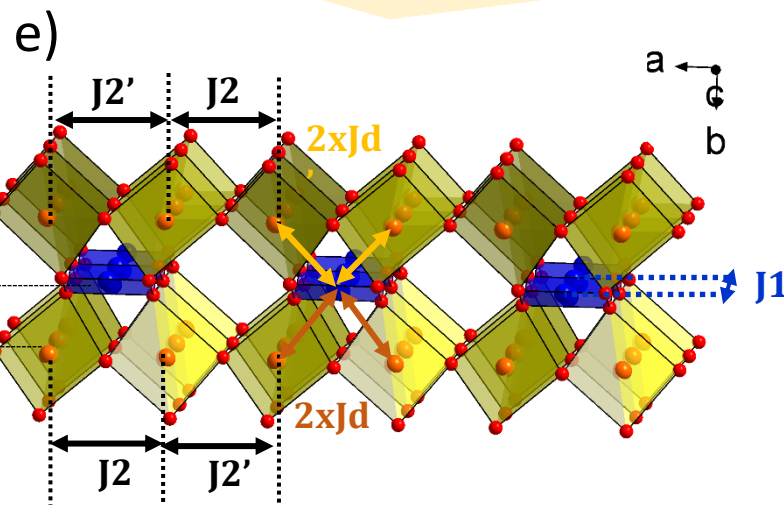
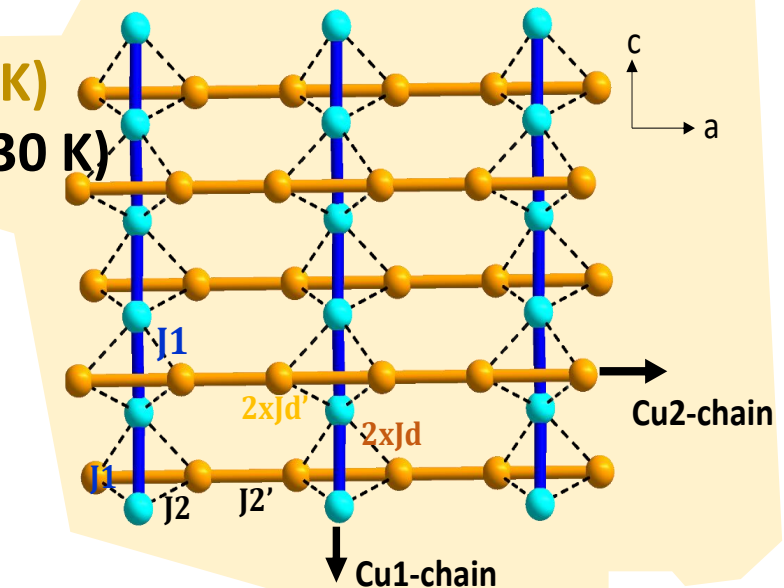
## • The $Cu_3Te_2O_5(OH)_4$ compound



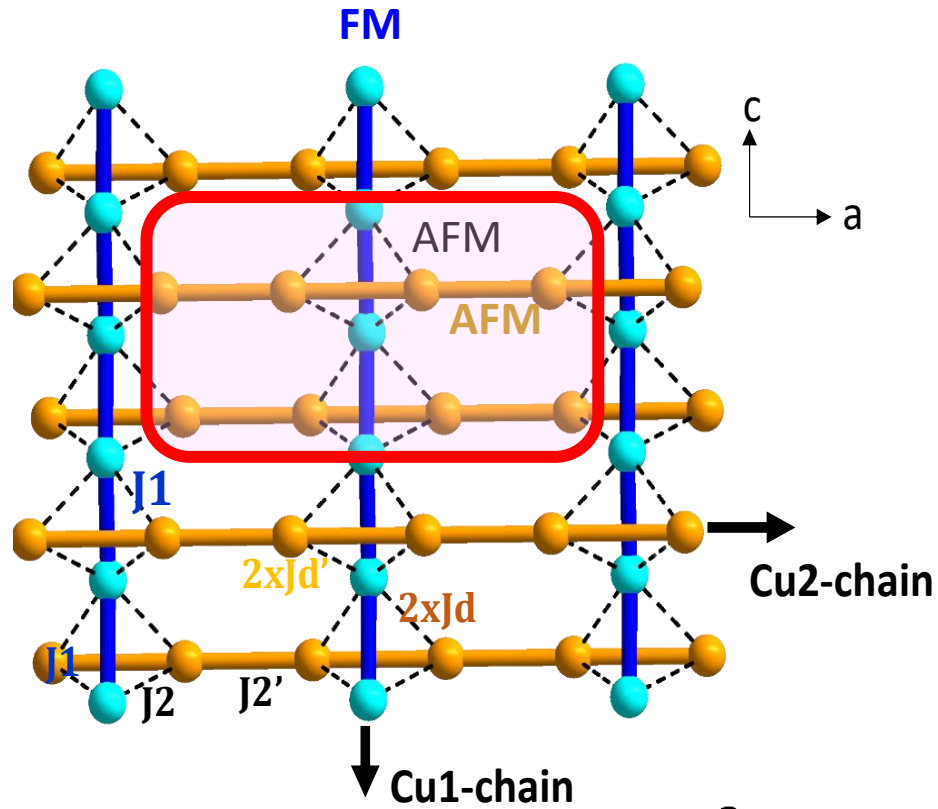
AFM :  $J_2/J_2'$  (=147.4/118.4 K)

AFM  $J_d$  (= 30 K)

FM :  $J_1$  (= -98 K)

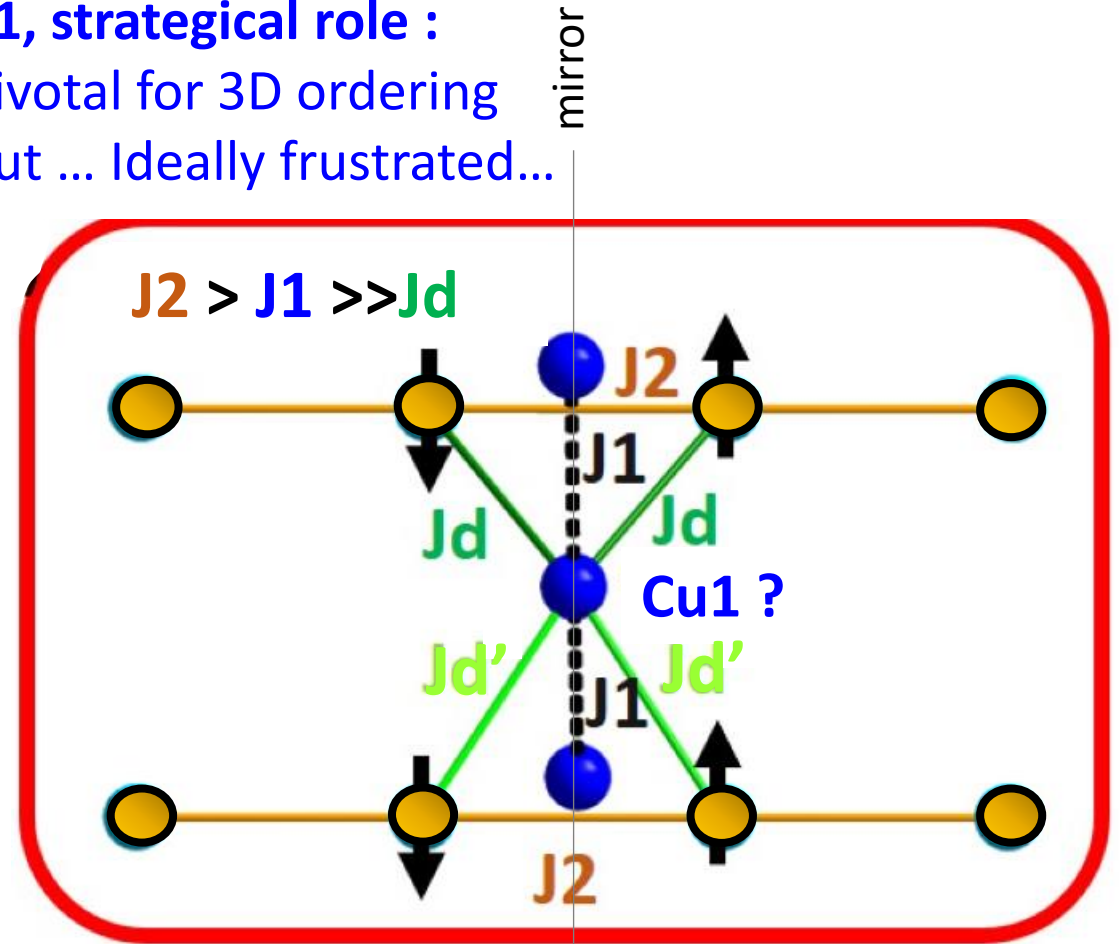


# Towards « true » Quantum Spin Liquid



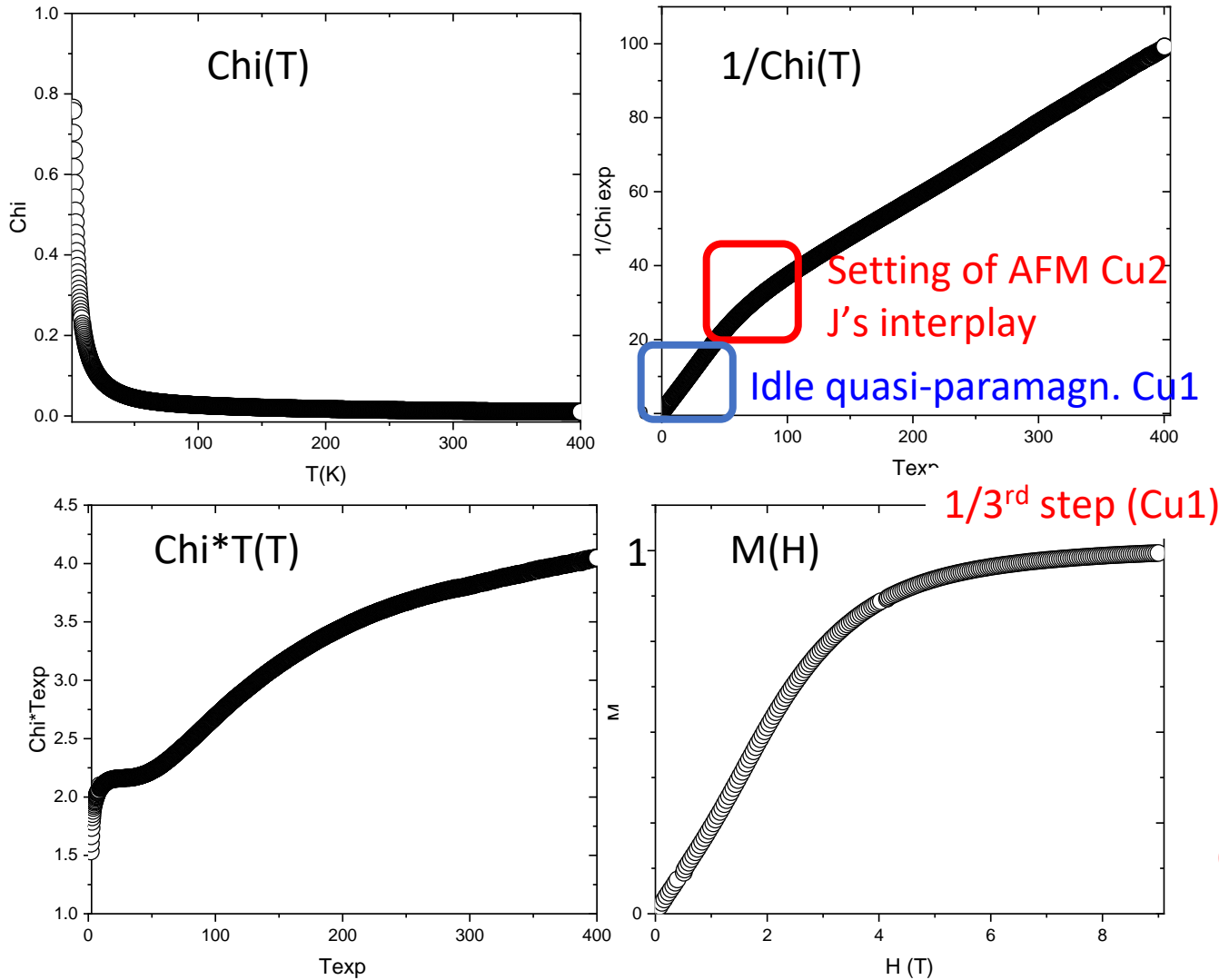
**Cu1, strategical role :**

- pivotal for 3D ordering
- but ... Ideally frustrated...



Cu1 ideally frustrated by the mirror-symmetry

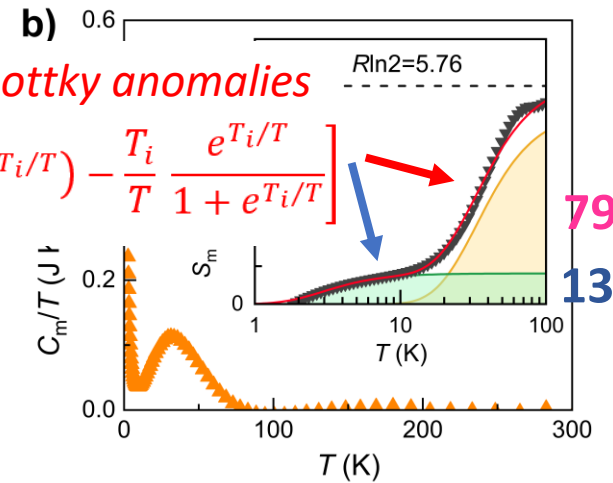
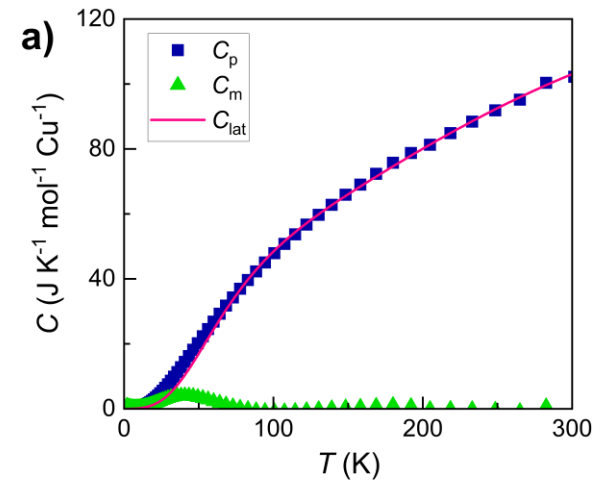
# Towards « true » Quantum Spin Liquid



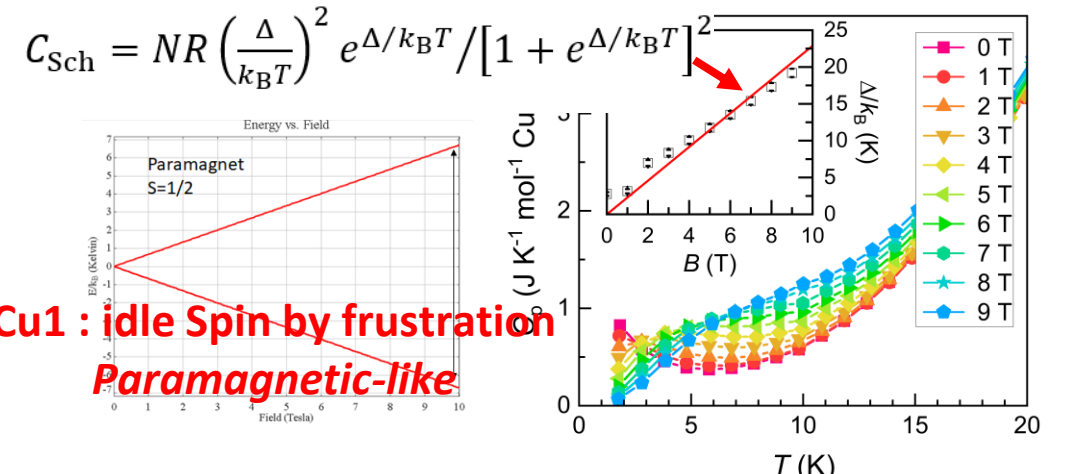
K.-Y. Choi  
 Center for Integrated Nanostructure Physics

$C_p \rightarrow$  two Schottky anomalies

$$S_m(T) = \sum_{i=1,2} a_i \left[ \ln(1 + e^{T_i/T}) - \frac{T_i}{T} \frac{e^{T_i/T}}{1 + e^{T_i/T}} \right]$$

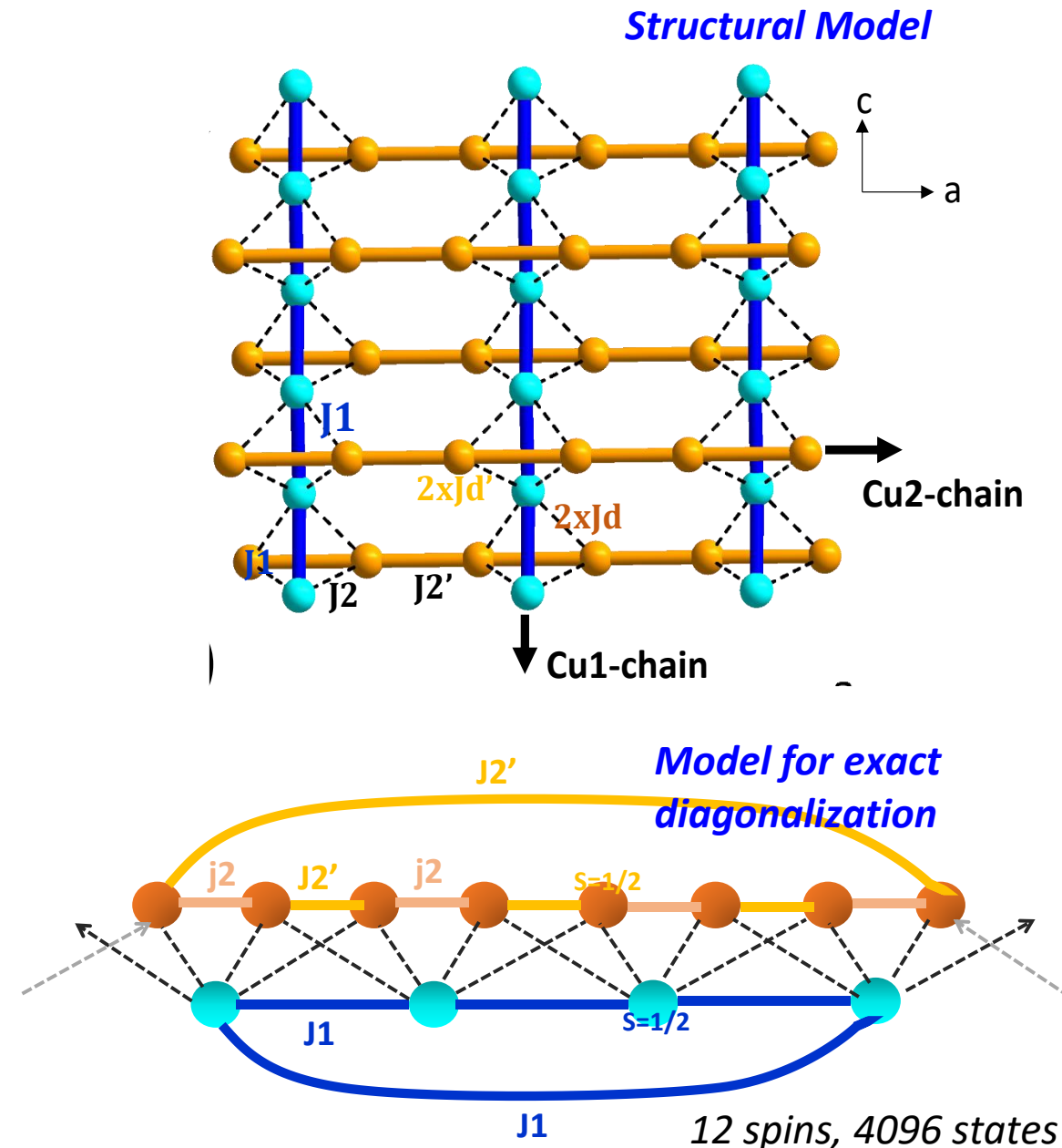
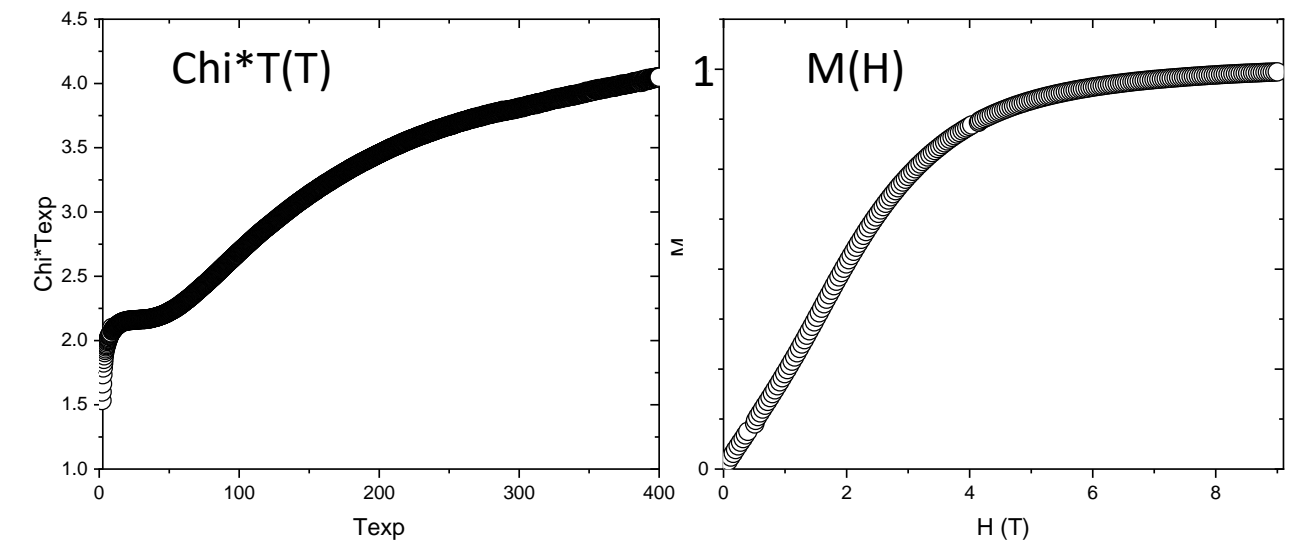
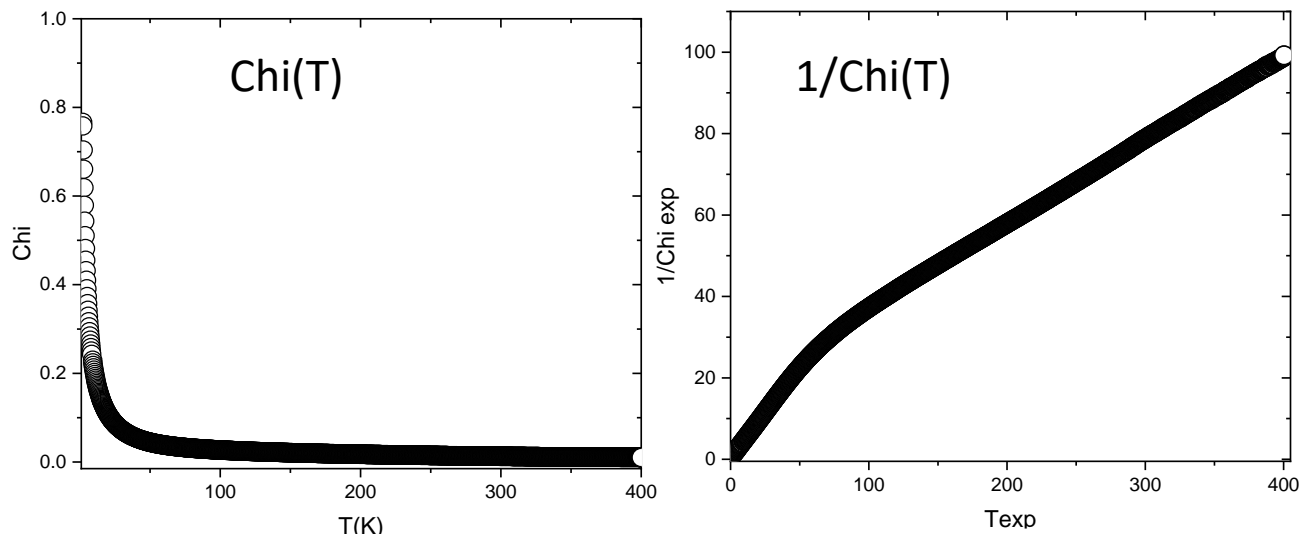


linear gap vs. H

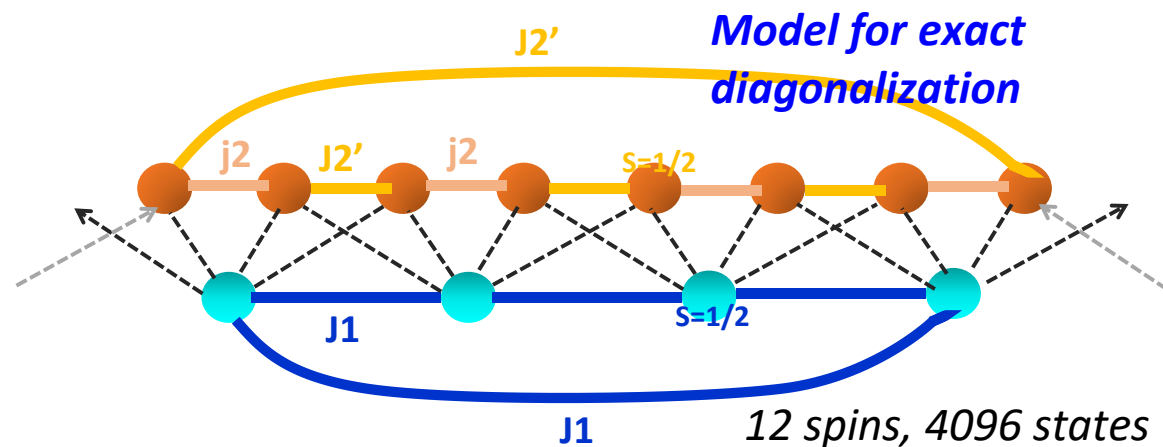
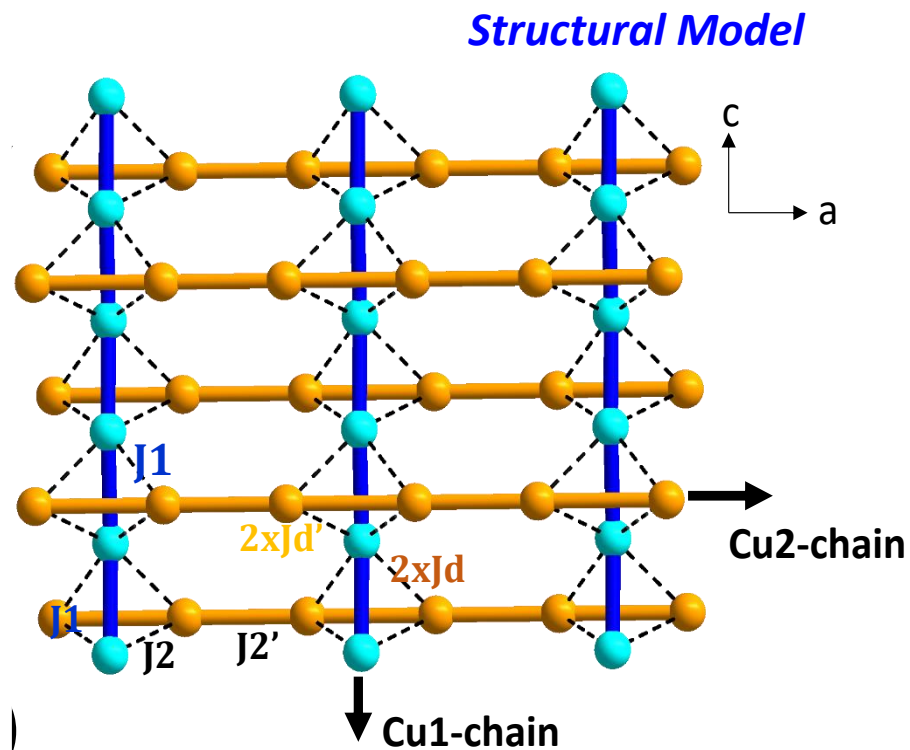
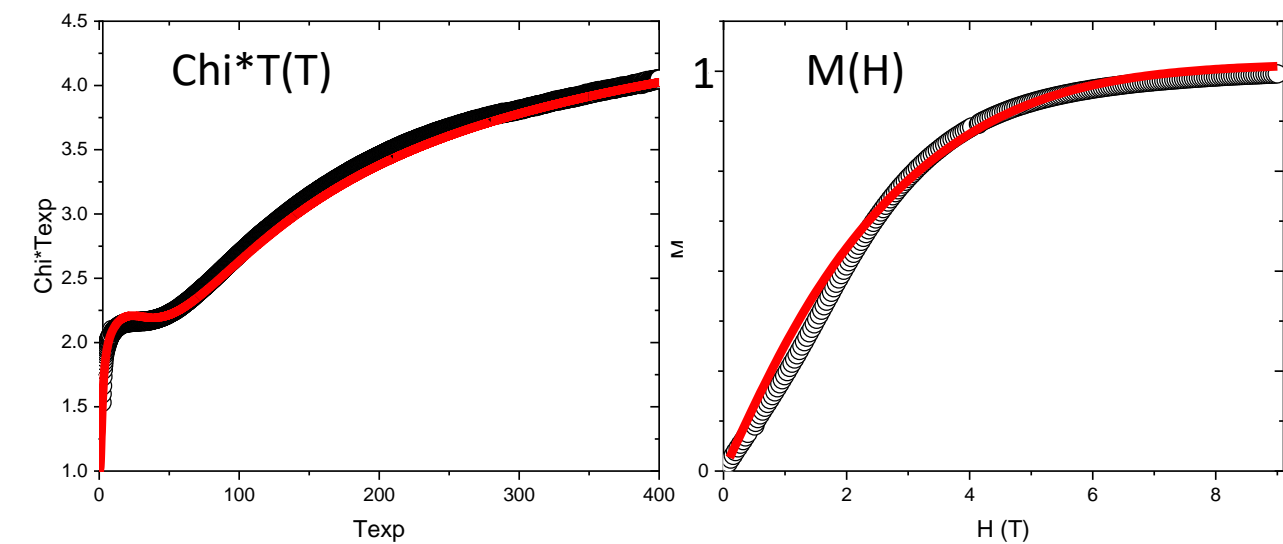
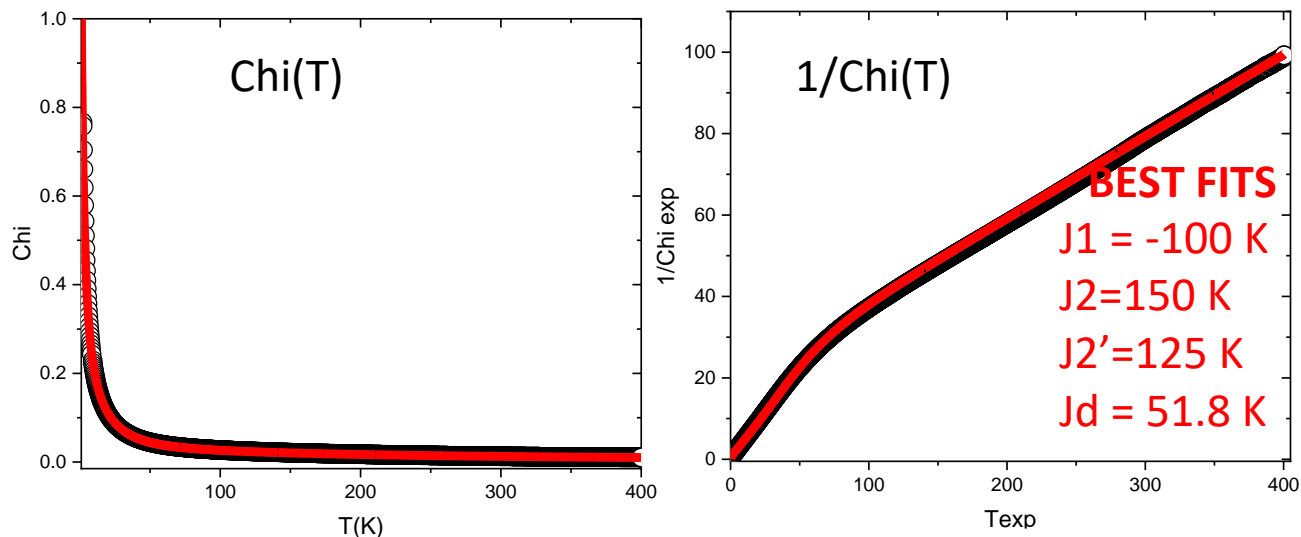


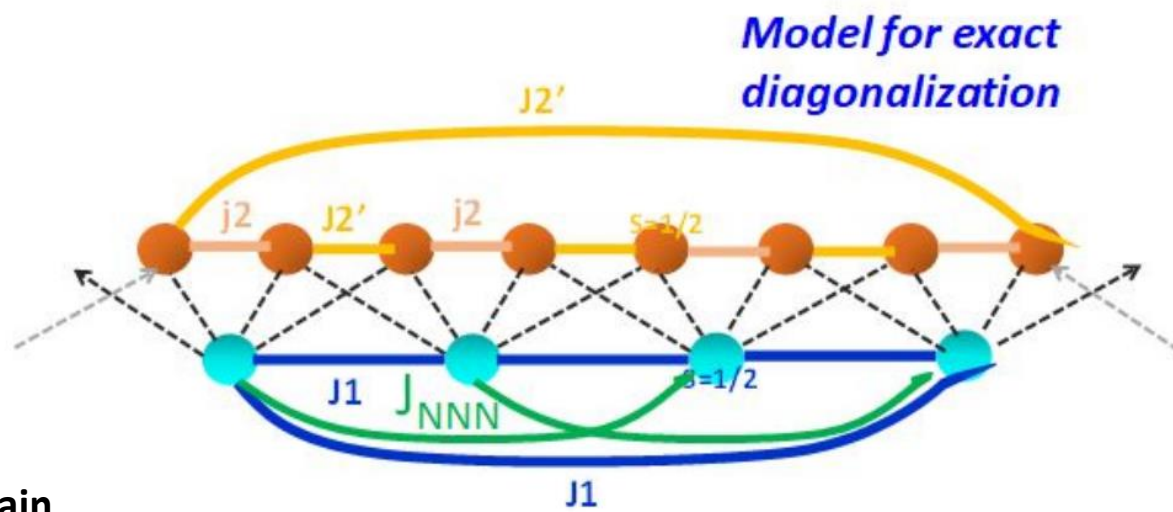
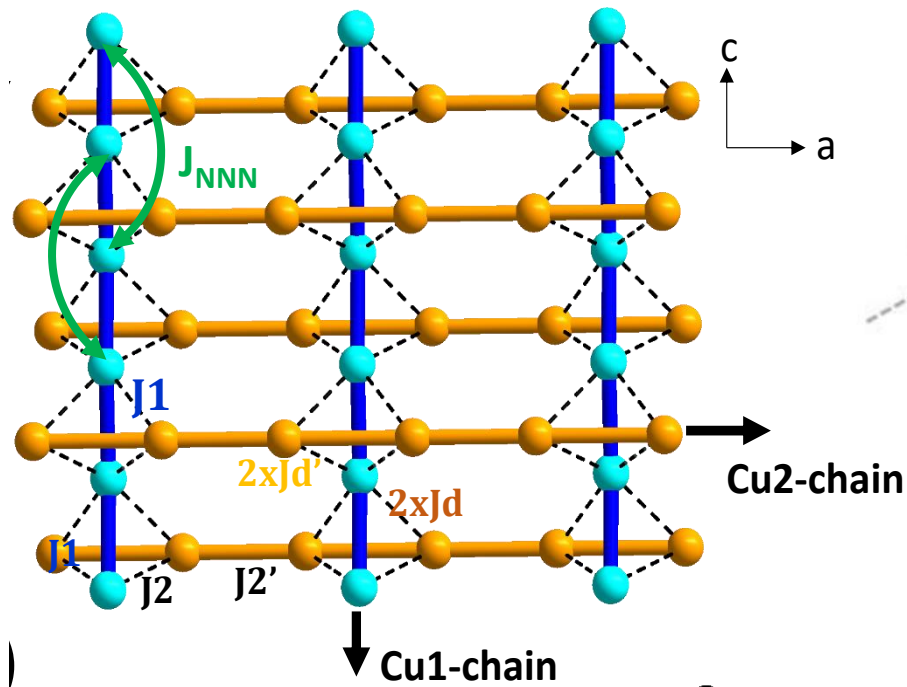
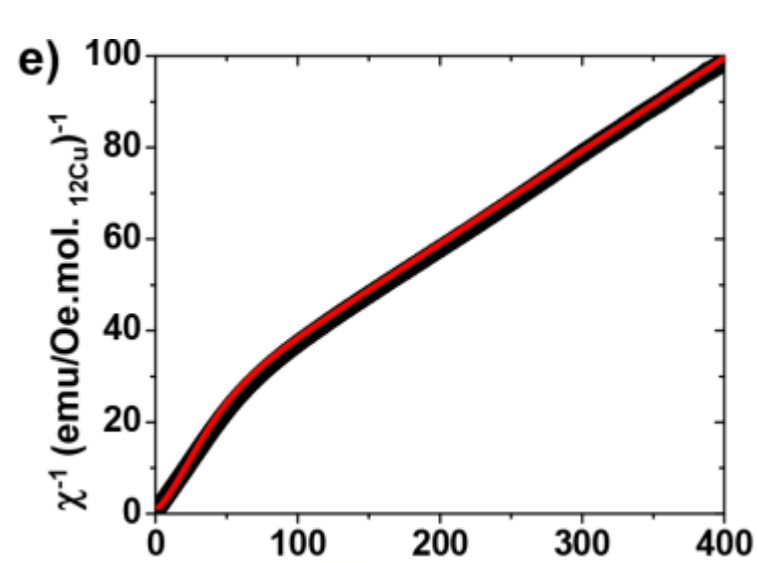
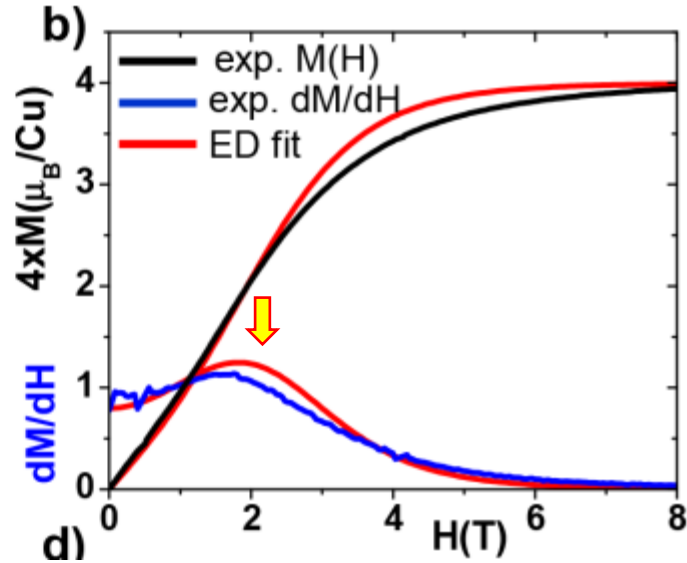
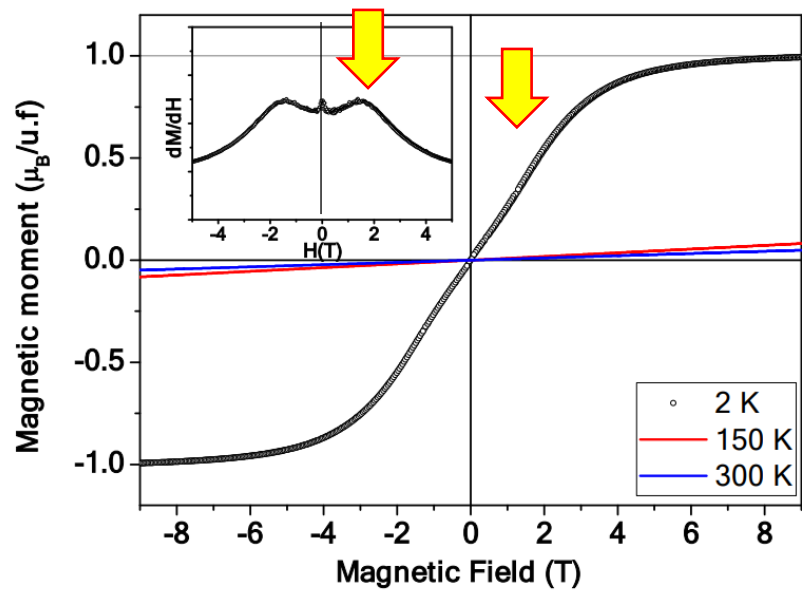
Cu1 : idle Spin by frustration  
 Paramagnetic-like

79% Cu2 + x% Cu1  
 13% x' %Cu1



## Quantum spin liquid





## $\text{Cu}_3\text{Te}_2\text{O}_5(\text{OH})_4$ : A Frustrated Two-Dimensional Quantum “Magnetic Raft” as a Possible Pathway to a Spin Liquid

Tianyu Zhu,<sup>▽</sup> Bei Zhu,<sup>▽</sup> Olivier Mentré,<sup>\*</sup> Suheon Lee, Dan Chen, Yanling Jin, Wenxuan Zhu, Angel M. Arévalo-López, Claire Minaud, Kwang-Yong Choi, and Minfeng Lü<sup>\*</sup>



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